# Tensor operations on characteristic functions of $C_{0}$ contractions 

H. BERCOVICI and D. VOICULESCU

By the results of [14], [15] and [1] every contraction $T$ of class $C_{0}$ acting on a separable Hilbert space is quasi-similar to a unique Jordan operator. If $T$ has finite defect indices then its Jordan model also shares this property and B. Sz.-Nagy and C. FoIAş proved in [14] that the determinant of the characteristic function of $T$ and of the Jordan model coincide in this case.

Also in the case of finite defect indices, from the work of E. A. Nordgren and B. Moore ([10] and [8]; cf. also [16]) it is known that the inner functions appearing in the Jordan model of $T$ can be computed from the minors of the determinant of the characteristic function of $T$.

It is an immediate problem to find characterizations for the inner functions in the Jordan model of a general $C_{0}$ contraction, and to look for special characterizations in the case of weak contractions of class $C_{0}$ ([13], chapter VIII) when the characteristic function has a determinant.

Also, the determinant being a representation of the unitary group on a finitedimensional space, more generally we may perform on the characteristic function of a contraction tensor operations of the type associated to irreducible representations of unitary groups, and ask about the properties of the operators having these functions as characteristic functions.

In the first part of this paper we consider tensor operations corresponding to irreducible representations of unitary groups applied to characteristic functions of operators of class $C_{0}$, the main result being that these operations preserve the quasisimilarity of the associated operators, provided the given operators have equal defect indices. This assertion is also adapted for the case of unequal defect indices, using impure characteristic functions.

As a corollary we characterize the inner functions in the Jordan model of a $C_{0}$ contraction by means of the smallest scalar inner multiples of the exterior

[^0]powers of the characteristic function. We also obtain estimates for the defect operator of a $C_{0}$ contraction in terms of the Jordan model.

In the second part of the paper we construct higher order algebraic adjoints of the characteristic function of a weak contraction. This enables us, using the results of the first part, to extend the above mentioned result of E. A. Nordgren and $B$. Moore to the case of weak contractions of class $C_{0}$.

We also prove that the determinant of the characteristic function of such a contraction is an inner function.

Using the results of the first part concerning defect operators, we prove that a $C_{0}$ contraction is a weak contraction, if and only if its Jordan model is a weak contraction. This extends a result of L. E. Isaev [5] on dissipative operators, which via Cayley transform (see [13] ch. IX) shows that a $C_{0}$ contraction with Jordan model $S\left(m_{a}\right), m_{a}(\lambda)=\exp (-a(1+\lambda) /(1-\lambda))(a>0)$, is a weak contraction.

## Part I

## § 1. Notation and preliminaries

1. We shall consider separable (finite or infinite dimensional) Hilbert spaces over the complex field $\mathbf{C}$.

We shall denote by $\mathfrak{H}, \boldsymbol{\Omega}, \ldots$ Hilbert spaces; $\langle.,$.$\rangle will denote the scalar product$ in any such space. If $\mathfrak{Y}$ is a subspace of $\mathfrak{G}$ we denote by $P_{\mathfrak{Y}}$ the orthogonal projection of $\mathfrak{5}$ onto $\mathfrak{Y}$ and by $\mathfrak{Y}^{\perp}$ or $\mathfrak{S} \ominus \mathfrak{Y}$ the orthogonal complement of $\mathfrak{Y}$. ( $\left.M\right)^{-}$denotes the norm-closure of the subset $M \subset \mathfrak{H}$. If $\left\{Y_{\alpha}\right\}_{\alpha \in A}$ is a family of subsets of $\mathfrak{H}, \bigvee_{\alpha \in A} Y_{\alpha}$ will denote the closed linear span of $\bigcup_{\alpha \in A} Y_{\alpha} . X \vee Y$ will denote the closed linear span of $X \cup Y$.

If $\mathfrak{S}$ and $\mathfrak{\Omega}$ are Hilbert spaces we shall denote by $\mathfrak{S} \otimes \mathfrak{A}$ their tensor product, which is also a Hilbert space. Recall that

$$
\begin{equation*}
\left\langle f \otimes g, f^{\prime} \otimes g^{\prime}\right\rangle=\left\langle f, f^{\prime}\right\rangle\left\langle g, g^{\prime}\right\rangle \text { for } f, g \in \mathfrak{G}, f^{\prime} g^{\prime} \in \mathfrak{S}^{\prime} \tag{1.1}
\end{equation*}
$$

$\mathfrak{S}^{\otimes n}$ will denote the tensor product $\mathfrak{S} \otimes \mathfrak{S} \otimes \ldots \otimes \mathfrak{S}$ ( $n$ times).
We denote by $\mathscr{L}(\mathfrak{G}, \mathfrak{\Omega})$ the linear space of all linear bounded operators $X: \mathfrak{S} \rightarrow \mathfrak{F}, \mathscr{L}(\mathfrak{G})=\mathscr{L}(\mathfrak{H}, \mathfrak{H})$. If $S$ is any subset of $\mathscr{L}(\mathfrak{H}),(S)^{\prime}$ denotes the commutant of $S . \mathscr{U}(\mathfrak{5})$ denotes the group of unitary operators on $\mathfrak{5}$.

If $T \in \mathscr{L}(\mathfrak{G})$, the operator $\Gamma_{n}(T) \in \mathscr{L}\left(\mathfrak{G}^{\otimes n}\right)$ is determined by
(1.2) $\quad \Gamma_{n}(T)\left(h_{1} \otimes h_{2} \otimes \ldots \otimes h_{n}\right)=T h_{1} \otimes T h_{2} \otimes \ldots \otimes T h_{n}, \quad h_{j} \in \mathcal{S} \quad(1 \leqq j \leqq n)$.

The map $\Gamma_{n}$ is multiplicative, commutes with the $*$-operation and restricted to $\mathscr{U}(\mathfrak{H})$ is a unitary representation.
2. Let us recall that $H^{\infty}$ is the Banach algebra of bounded analytic functions in the unit disc $D=\{z \in \mathbf{C}| | z \mid<1\}$. We denote by $H_{i}^{\infty}$ the set of inner functions in $H^{\infty}$, that is $m \in H_{i}^{\infty}$ if and only if $m$ has (dt-)almost everywhere radial limits $m\left(e^{i t}\right)$ of modulus one. We shall abuse notation sometimes, writing $m=m^{\prime}$ for two inner functions such that $\mathrm{m} / \mathrm{m}^{\prime}$ is a constant (of modulus one).

If $\left\{f_{\alpha}\right\}_{\alpha \in A}$ is a family of $H^{\infty}$-functions, not all 0 , we denote by $\bigwedge_{\alpha \in A} f_{\alpha}$ the greatest common inner divisor of the functions $f_{\alpha}$.

Consider also the Hardy space $H^{2}$ and, for a Hilbert space $\mathfrak{5}$, the vector-valued Hardy space $H^{2}(\mathfrak{H})$ which can be identified with $\mathfrak{G} \otimes H^{2}$.

If $T \in \mathscr{L}(\mathfrak{H})$ and $S \in \mathscr{L}\left(H^{2}\right)$ we shall consider $T \otimes S$ as an operator on $H^{2}(\mathfrak{H})$. For $f \in H^{\infty}(\mathfrak{H}), g \in H^{\infty}(\Omega)$ we shall denote (somewhat ambiguously) by $f \otimes g$ the element of $H^{2}(\mathfrak{G} \otimes \mathfrak{R})$ defined by

$$
\begin{equation*}
(f \otimes g)(z)=f(z) \otimes g(z), \quad z \in D \tag{1.3}
\end{equation*}
$$

For any two Hilbert spaces $\mathfrak{5}, \mathcal{R}$ the operator-valued Hardy space $H^{\infty}(\mathscr{L}(\mathfrak{H}, \mathfrak{R}))$ is the set of all bounded, $\mathscr{L}(\mathfrak{F}, \mathfrak{\Omega})$-valued analytic functions in the unit disc.

A function $\Theta \in H^{\infty}(\mathscr{L}(\mathfrak{H}, \Omega))$ is contractive if $\|\Theta(z)\| \leqq 1, z \in D$. Any function $\Theta \in H^{\infty}(\mathscr{L}(\mathfrak{H}, \mathfrak{\Omega}))$ may be considered as an element of $\mathscr{L}\left(H^{2}(\mathfrak{H}), H^{2}(\mathfrak{\Omega})\right)$ that commutes with scalar $H^{\infty}$-multiplications.

We say that two functions

$$
\Theta_{i} \in H^{\infty}\left(\mathscr{L}\left(\mathfrak{S}_{i}, \mathfrak{\Re}_{i}\right)\right) \quad(i=1,2)
$$

coincide if there are unitary operators $U: \mathfrak{S}_{1} \rightarrow \mathfrak{S}_{2}, V: \mathfrak{\Re}_{1} \rightarrow \mathfrak{R}_{2}$ such that $\Theta_{2}(\lambda) U=$ $=V \Theta_{1}(\lambda)$ for all $\lambda \in D$.

A function $\Theta \in H^{\infty}(\mathscr{L}(\mathfrak{H}, \boldsymbol{\Omega}))$ is inner if it is isometric as an element of $\mathscr{L}\left(H^{2}(\mathfrak{H}), H^{2}(\Omega)\right) . \Theta$ is *-inner if the function $\Theta^{\sim}$ defined by

$$
\begin{equation*}
\Theta^{\sim}(z)=\Theta(\bar{z})^{*}, \quad z \in D \tag{1.4}
\end{equation*}
$$

is inner. $\Theta$ is two-sided inner if it is simultaneously inner and $*$-inner. We denote by $H_{i}^{\infty}(\mathscr{L}(\mathfrak{S}, \mathfrak{\Omega}))$ the set of two-sided inner functions in $H^{\infty}(\mathscr{L}(\mathfrak{H}, \mathfrak{\Omega}))$.

For any $\Theta \in H^{\infty}(\mathscr{L}(\mathfrak{H}))$ we denote by $\Gamma_{n}(\Theta)$ the element of $H^{\infty}\left(\mathscr{L}\left(\mathfrak{G}^{\otimes n}\right)\right)$ defined by

$$
\begin{equation*}
\left(\Gamma_{n}(\Theta)\right)(z)=\Gamma_{n}(\Theta(z)), \quad z \in D \tag{1.5}
\end{equation*}
$$

If $\Theta \in H_{i}^{\infty}(\mathscr{L}(\mathfrak{H}))$ then $\Gamma_{n}(\Theta) \in H_{i}^{\infty}\left(\mathscr{L}\left(\mathfrak{S}^{\otimes n}\right)\right)$.
3. For any $\Theta \in H_{i}^{\infty}(\mathscr{L}(\mathfrak{H}))$ we define $S(\Theta)$ as the operator acting on

$$
\begin{equation*}
\mathfrak{H}(\Theta)=H^{2}(\mathfrak{H}) \ominus \Theta H^{2}(\mathfrak{H}) \tag{1.6}
\end{equation*}
$$

and defined by

$$
\begin{equation*}
\left(S(\Theta)^{*} u\right)(z)=z^{-1}(u(z)-u(0)), \quad z \in D, \quad u \in \mathfrak{H}(\Theta) \tag{1.7}
\end{equation*}
$$

If $\Theta$ is pure then it coincides with the characteristic function of $S(\Theta)$ and in this case $\operatorname{dim} \mathfrak{H}$ equals the defect indices of $S(\Theta)$ [13]. Recall that, for a contraction $T \in \mathscr{L}(\Omega)$, the defect operators are $D_{T}=\left(I-T^{*} T\right)^{1 / 2}, D_{T^{*}}=\left(1-T T^{*}\right)^{1 / 2}$ and the defect indices $\mathrm{D}_{T}, \mathrm{D}_{T^{*}}$ are the ranks of $D_{T}$ and $D_{T^{*}}$, respectively.

Let $\mu_{T}$ denote the multiplicity of $T$, i.e. the least cardinal of cyclic sets for $T$.
We shall need the lifting of commutants theorem of [13] in the following form. If $\Theta \in H_{i}^{\infty}(\mathscr{L}(\mathfrak{H})), \Theta^{\prime} \in H_{i}^{\infty}\left(\mathscr{L}\left(\mathfrak{5}^{\prime}\right)\right)$ and $X \in \mathscr{L}\left(\mathfrak{H}(\Theta), \mathfrak{H}\left(\Theta^{\prime}\right)\right)$ satisfy the relation

$$
S\left(\Theta^{\prime}\right) X=X S(\Theta)
$$

then there is an $A \in H^{\infty}\left(\mathscr{L}\left(5, \mathfrak{G}^{\prime}\right)\right)$ such that

$$
\begin{gather*}
A \Theta H^{2}(\mathfrak{H}) \subset \Theta^{\prime} H^{2}\left(\mathfrak{H}^{\prime}\right) \quad \text { and }  \tag{1.8}\\
X h=P_{\mathfrak{5}\left(\Theta^{\prime}\right)} A h, \quad h \in \mathfrak{H}(\Theta) . \tag{1.9}
\end{gather*}
$$

The operator $X$ is one-to-one if and only if, for $h \in H^{2}(\mathfrak{H})$,

$$
\begin{equation*}
h \in \Theta H^{2}(\mathfrak{H}) \Leftrightarrow A h \in \Theta^{\prime} H^{2}\left(\mathfrak{H}^{\prime}\right), \tag{1.10}
\end{equation*}
$$

and has dense range if and only if

$$
\begin{equation*}
A H^{2}(\mathfrak{H}) \vee \Theta^{\prime} H^{2}\left(\mathfrak{H}^{\prime}\right)=H^{2}\left(\mathfrak{H}^{\prime}\right) \tag{1.11}
\end{equation*}
$$

Let us recall that $X$ is called a quasi-affinity if it is one-to-one and has dense range.
The operator $S(\Theta)$ is of class $C_{0}$ if and only if $\Theta$ has a scalar multiple, that is, if

$$
\begin{equation*}
\Theta H^{2}(\mathfrak{H}) \supset m H^{2}(\mathfrak{H}) \tag{1.12}
\end{equation*}
$$

for some $m \in H_{i}^{\infty}$. The minimal function of $T=S(\Theta)$ is then the greatest common inner divisor $m_{T}$ of the functions $m$ satisfying (1.12) [13].

A Jordan operator is an operator $S(\Theta)$ determined by a function of the form

$$
\Theta=\left[\begin{array}{lllll}
m_{1} & & & 0 & \\
& m_{2} & & \\
& & \cdot & \\
0 & & & \\
& & & & \cdot
\end{array}\right]
$$

where $m_{j} \in H_{i}^{\infty}$ and $m_{j+1}$ divides $m_{j}$ for each $j$. We shall denote it also by $S(M)$, $M=\left\{m_{j}\right\}_{j=1}^{\infty}$. By the results of [14], [15], [1] every $C_{0}$ contraction acting on a separable Hilbert space is quasisimilar to a unique Jordan model $S(M)$.
4. For a finite group $G$ we shall denote by $C^{*}(G)$ the $C^{*}$-algebra of $G$ [2], and by $\hat{G}$ the set of all (equivalence classes) of irreducible unitary representations of $G$. The elements of $C^{*}(G)$ will be written in the form $\sum_{g \in G} c_{g} g$ where $c_{g} \in \mathbf{C}$, so that for any unitary representation $\pi$ of $G$ the associated representation of $C^{*}(G)$ is
given by

$$
\pi\left(\sum_{g \in G} c_{g} g\right)=\sum_{g \in G} c_{g} \pi(g)
$$

Let $\mathfrak{S}_{n}$ be the group of permutations of the set $\{1,2, \ldots, n\}$. The group $\mathfrak{S}_{n-1}$ will be identified with the subgroup of $\mathbb{S}_{n}$ consisting of those permutations of $\mathbb{S}_{n}$ that leave $n$ fixed and $C^{*}\left(\mathfrak{S}_{n-1}\right)$ will be considered as a sub-algebra of $C^{*}\left(\mathfrak{S}_{n}\right)$.
$\hat{\Theta}_{n}$ is known to be in one-to-one correspondence with signatures $\tau=\left(t_{1} \geqq \ldots \geqq t_{n}\right)$, $t_{j}$ non-negative integers, $\sum_{j=1}^{n} t_{j}=n$, and the corresponding minimal central projections $p_{\tau}$ of $C^{*}\left(\Theta_{n}\right)$ are given by the central Youngsymmetrizers [18], [6], [9]. It is known [17], Ch. $\mathrm{V}, \S 18$, that an irreducible representation of signature $\tau=\left(t_{1} \geqq t_{2} \geqq \ldots \geqq t_{n}\right)$ restricted to $\Im_{n-1}$ contains the irreducible representation of signature $\tau^{\prime}=\left(t_{1}^{\prime} \geqq t_{2}^{\prime} \geqq \ldots\right.$ $\ldots \geqq t_{n-1}^{\prime}$ ) if and only if

$$
\begin{equation*}
t_{1} \geqq t_{1}^{\prime} \geqq t_{2} \geqq t_{2}^{\prime} \geqq \ldots \geqq t_{n-1} \geqq t_{n-1}^{\prime} \geqq t_{n} \tag{1.13}
\end{equation*}
$$

(this will be written $\tau^{\prime}<\tau$ ) and that the multiplicity of $\tau^{\prime}$ is one in this case.
Consider now a Hilbert space $\mathcal{K}$. On $\boldsymbol{\Omega}^{\otimes n}$ there is a unitary representation $\pi_{n}$ of $\widehat{\Xi}_{n}$ given by

$$
\begin{equation*}
\pi_{n}(\sigma)\left(k_{1} \otimes \ldots \otimes k_{n}\right)=k_{\sigma^{-1}(1)} \otimes \ldots \otimes k_{\sigma^{-1}(n)}, \quad \sigma \in \mathbb{S}_{n} \tag{1.14}
\end{equation*}
$$

By one of the basic results of Hermann Weyl ([18], [6], see also [11], [7] for the adaptation to the case when $\operatorname{dim} \Omega$ is infinite) we have

$$
\begin{equation*}
\left(\Gamma_{n}(\mathscr{U}(\mathscr{\Omega}))\right)^{\prime}=\left(\Gamma_{n}(\mathscr{L}(\mathscr{\Omega}))\right)^{\prime}=\pi_{n}\left(C^{*}\left(\mathscr{S}_{n}\right)\right) \tag{1.15}
\end{equation*}
$$

The irreducible representations of $\mathscr{U}(\boldsymbol{\Omega})$ which will be considered are also labelled by signatures, so we shall first make a convention. A signature will be a decreasing sequence $\tau=\left(t_{1} \geqq t_{2} \geqq \ldots\right)$ of nonnegative integers, of finite or infinite length $l(\tau)$. By $\downarrow(\tau)$ we shall denote the number of nonzero elements among the $t_{j}$ 's and $|\tau|$ will stand for $\sum_{j=1}^{l(\tau)} t_{j}$.

Thus for instance the set $\hat{\Xi}_{n}$ is in a one-to-one correspondence with those signatures $\tau$ for which $l(\tau)=|\tau|=n$. Two signatures $\tau=\left(t_{1} \geqq t_{2} \geqq \ldots\right)$ and $\tau^{\prime}=$ $=\left(t_{1}^{\prime} \geqq t_{2}^{\prime} \geqq \ldots\right)$ are essentially equivalent if $\_(\tau)={ }_{\perp}\left(\tau^{\prime}\right)$ and $t_{j}=t_{j}^{\prime}$ for $j=1,2, \ldots, \perp(\tau)$.

For a signature $\tau$ with $l(\tau)=\operatorname{dim} \Omega,|\tau|<\infty$, there corresponds an irreducible representation $\varrho_{\tau}$ of $\mathscr{U}(\Omega)$ on a Hilbert space $\boldsymbol{\Omega}^{\tau}$ (these are the irreducible representations of "positive" signatures; cf. [18], [6] for the case $\operatorname{dim} \Omega<\infty$ and [11] for the extension to the case $\operatorname{dim} \Omega=\infty$ ).

The representation $\varrho_{\tau}$ can be defined as follows: consider $\tilde{\tau}$, the signature of length $|\tau|$ essentially equivalent to $\tau$, and let $q_{\tilde{\tau}}$ be any minimal projection in $C^{*}\left(\mathcal{G}_{|\tau|}\right)$ such that $q_{\tilde{z}} \leqq p_{\tilde{\imath}}$. Then $\varrho_{\tau}$ is defined as the restriction of $\Gamma_{|\tau|}$ to $\pi_{|\tau|}\left(q_{\tilde{z}}\right) \mathcal{S}^{\otimes|\tau|}$. Clearly
$\varrho_{\tau}$ extends to a multiplicative homomorphism of the multiplicative semigroup $\mathscr{L}(\Omega)$ which is holomorphic. Also clearly the restriction of $\Gamma_{|\tau|}$ to $\pi_{|\tau|}\left(p_{\bar{\tau}}\right) \Omega^{\otimes|r|}$ is a finite multiple of $\varrho_{\tau}$.

Another classical fact we need is that for $\tau$ with $l(\tau)=|\tau|=n$ we have $\pi_{n}\left(p_{\imath}\right) \neq 0$ if and only if $\Omega(\tau) \leqq \operatorname{dim} \boldsymbol{\Omega}$.

## § 2. Tensor operations on operator-valued functions

Let $\Omega$ be a Hilbert space. For any $k \in \Omega$ we shall consider the map $T_{k}: \Omega^{\otimes n} \rightarrow$ $\rightarrow \Omega^{\otimes(n+1)}$ defined by

$$
\begin{equation*}
T_{k}\left(k_{1} \otimes k_{2} \otimes \ldots \otimes k_{n}\right)=k_{1} \otimes \ldots \otimes k_{n} \otimes k \tag{2.1}
\end{equation*}
$$

Clearly $T_{k}$ is proportional to an isometry and

$$
\begin{equation*}
T_{k}^{*}\left(k_{1} \otimes \ldots \otimes k_{n+1}\right)=\left\langle k_{n+1}, k\right\rangle k_{1} \otimes \ldots \otimes k_{n} \tag{2.2}
\end{equation*}
$$

Lemma 2.1. Consider two signatures $\tau^{\prime}<\tau, \quad l\left(\tau^{\prime}\right)=\left|\tau^{\prime}\right|=n, \quad l(\tau)=|\tau|=n+1$ such that $\&(\tau) \leqq \operatorname{dim} \AA$. Then we have:

$$
\begin{equation*}
\underset{k \in \Omega}{\bigvee} \pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right) \quad \Omega^{\otimes(n+1)}=\pi_{n}\left(p_{\tau^{\prime}}\right) \Omega^{\otimes n} \tag{2.3}
\end{equation*}
$$

Proof. Let us denote by $\mathfrak{F}$ the space on the left hand side of (2.3). Then $\mathcal{F}$ is $\pi_{n}\left(\mathbb{S}_{n}\right)$-invariant and $\Gamma_{n}(\mathscr{U}(\Omega))$-invariant.

Indeed, for $\sigma \in \Im_{n}$ we have

$$
\pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right) \pi_{n+1}(\sigma)=\pi_{n}(\sigma) \pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right)
$$

since $p_{\tau^{\prime}}, p_{\tau}$ commute with $C^{*}\left(\Theta_{n}\right)$ and $T_{k} \pi_{n}(\sigma)=\pi_{n+1}(\sigma) T_{k}$. Also,

$$
\mathfrak{\Re}^{\otimes n} \ominus \mathscr{F}=\bigcap_{k \in \mathcal{A}} \operatorname{Ker}\left[\pi_{n+1}\left(p_{\imath}\right) T_{k} \pi_{n}\left(p_{\tau^{\prime}}\right)\right]
$$

and for any $U \in \mathscr{U}(\boldsymbol{\Omega})$ we have

$$
\Gamma_{n}(U) \operatorname{Ker}\left[\pi_{n+1}\left(p_{\tau}\right) T_{k} \pi_{n}\left(p_{\tau^{\prime}}\right)\right]=\operatorname{Ker}\left[\pi_{n+1}\left(p_{\tau}\right) T_{U k} \pi_{n}\left(p_{\tau^{\prime}}\right)\right]
$$

so that $\mathcal{\Omega}^{\otimes n} \ominus \mathscr{F}$ is invariant for $\Gamma_{n}(\mathscr{U}(\mathcal{\Omega}))$ and hence so is $\mathscr{F}$.
Therefore $P_{\mathfrak{F}} \in\left(\pi_{n}\left(C^{*}\left(\mathcal{S}_{n}\right)\right) \cup \Gamma_{n}(\mathscr{U}(\mathcal{N}))\right)^{\prime}$ and $P_{\mathfrak{F}} \leqq \pi_{n}\left(p_{\tau^{\prime}}\right)$. Hence by Hermann Weyl's theorem and because of the minimality of $p_{\tau^{\prime}}$ in the center of $C^{*}\left(\Theta_{n}\right)$ either $P_{\mathfrak{J}}=0$ or $P_{\mathfrak{J}}=\pi_{n}\left(p_{\tau^{\prime}}\right)$. So it will be sufficient to prove that $\mathfrak{F} \neq\{0\}$.

Observe that $\pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right)=T_{k}^{*} \pi_{n+1}\left(p_{\tau^{\prime}} p_{\tau}\right)$. On the other hand, $p_{\tau}$ is the central support of $p_{\tau} p_{\tau}$ in $C^{*}\left(\Im_{n+1}\right)$ as explained in the next paragraph. Thus, from $\pi_{n+1}\left(p_{\tau}\right) \neq 0$ we infer $\pi_{n+1}\left(p_{\tau^{\prime}} p_{\tau}\right) \neq 0$. Now $\bigcap_{k \in K} \operatorname{Ker} T_{k}^{*}=\{0\}$ so we can find $k \in \Omega$ such that $T_{k}^{*} \pi_{n+1}\left(p_{\tau} p_{\tau}\right) \neq 0$.

If $\varrho$ is an irreducible representation of the finite-dimensional $C^{*}$-algebra $A$, there is a minimal central projection $p$ of $A$ such that $\operatorname{ker} \varrho=(1-p) A$. Let $A_{1} \subset A_{2}$ be finite dimensional $C^{*}$-algebras with $1_{A_{2}} \in A_{1}, \varrho_{i}$ irreducible representations of $A_{i}$, and $p_{i}$ the corresponding minimal central projection of $A_{i}(i=1,2)$. Then $\varrho_{2} \mid A_{1}$ contains $\varrho_{1}$ if and only if $p_{1} p_{2} \neq 0$. Indeed, if $\varrho_{2} \mid A_{1}$ contains $\varrho_{1}$ we obviously have $\operatorname{ker}\left(\varrho_{2} \mid A_{1}\right) \subset \operatorname{ker} \varrho_{1}$, so that $p_{1} p_{2} \neq 0$ (since $p_{1} \notin \operatorname{ker} \varrho_{1}$ ). Conversely, if $p_{1} p_{2} \neq 0$ the two-sided ideal $J=\left\{x \in A_{1} ; p_{1} p_{2} x=0\right\}$ of $A_{1}$ contains ker $\varrho_{1}$ and $p_{1} \notin J$. Since $\varrho_{1}$ is irreducible and $A_{1}$ is finite-dimensional, ker $\varrho_{1}$ is a maximal ideal of $A_{1}$, so that $J=\operatorname{ker} \varrho_{1}$. It follows that $\operatorname{ker}\left(\varrho_{2} \mid A_{1}\right) \subset$ ker $\varrho_{1}$ and this in turn implies that $\varrho_{2} \mid A_{1}$ contains $\varrho_{1}$.

This completes the proof.
Lemma 2.2. Consider two signatures $\tau^{\prime}<\tau, l\left(\tau^{\prime}\right)=\left|\tau^{\prime}\right|=n, l(\tau)=|\tau|=n+1$, such that $s(\tau) \leqq \operatorname{dim} \Omega$ and let $\Theta \in H^{\infty}(\mathscr{L}(\Omega))$. For any $k \in \Omega$ we have:

$$
\begin{gather*}
\left(\left(\pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right)\right) \otimes I_{H^{2}}\right) \Gamma_{n+1}(\Theta) H^{2}\left(\Omega^{\otimes(n+1)}\right) \subset  \tag{2.4}\\
\subset\left(\left(\pi_{n}\left(p_{\tau^{\prime}}\right) \otimes I_{H^{2}}\right) \Gamma_{n}(\Theta) H^{2}\left(\Omega^{\otimes n}\right)\right)-
\end{gather*}
$$

Proof. Clearly both terms of (2.4) are invariant with respect to multiplication operators by scalar $H^{\infty}$-functions. Hence it is easily seen that it will be enough to prove that a function of the form

$$
z \rightarrow \pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right)\left(\Theta(z) k_{1} \otimes \ldots \otimes \Theta(z) k_{n+1}\right)
$$

is in

$$
\left(\pi_{n}\left(p_{\tau^{\prime}}\right) \otimes I_{H^{2}}\right) \Gamma_{n}(\Theta) H^{2}\left(\Omega^{\otimes n}\right)
$$

Writing $p_{\tau}=\sum_{\sigma \in \mathbb{S}_{n+1}} c_{\sigma} \sigma$ the assertion becomes obvious from the following computation:

$$
\begin{gathered}
\pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\tau}\right)\left(\Theta(z) k_{1} \otimes \ldots \otimes \Theta(z) k_{n+1}\right)= \\
=\pi_{n}\left(p_{\tau^{\prime}}\right) T_{k}^{*} \sum_{\sigma \in \mathbb{S}_{n+1}} c_{\sigma}\left(\Theta(z) k_{\sigma^{-1}(1)} \otimes \ldots \otimes \Theta(z) k_{\sigma^{-1}(n+1)}\right)= \\
=\sum_{\sigma \in \mathbb{S}_{n+1}} c_{\sigma}\left\langle\Theta(z) k_{\sigma^{-1}(n+1)}, k\right\rangle \pi_{n}\left(p_{\tau^{\prime}}\right) \Gamma_{n}(\Theta(z))\left(k_{\sigma^{-1}(1)} \otimes \ldots \otimes k_{\sigma^{-1}(n)}\right) .
\end{gathered}
$$

Let us now consider $\Theta \in H^{\infty}(\mathscr{L}(\boldsymbol{\Omega}))$ and let $\tau$ be a signature with $|\tau|<\infty$ and $\iota(\tau)=\operatorname{dim} \Omega$. Consider also $\tilde{\tau}$, the signature of length $|\tau|$ essentially equivalent to $\tau$. We define an inner function $d^{\tau}(\Theta)$ by

$$
\begin{equation*}
d^{\tau}(\Theta)=\wedge\left\{m \in H_{i}^{\infty} \mid m H^{2}\left(\Omega^{\tau}\right) \subset\left(\varrho_{\tau}(\Theta) H^{2}\left(\Omega^{\tau}\right)\right)^{-}\right\} \tag{2.5}
\end{equation*}
$$

(by convention we put $\wedge \varnothing=0, \varnothing$-the empty set).
Remark that in case $\Theta$ is an inner function, $\varrho_{\tau}(\Theta)$ is still an inner function and $d^{\tau}(\Theta)$ is the minimal function of $S\left(\varrho_{\tau}(\Theta)\right)$ in case $\varrho_{\tau}(\Theta)$ has a scalar multiple and zero otherwise. In case $\tau$ is of the form ( $1,1, \ldots, 1,0, \ldots$ ) with $j$ nonzero terms,
that is, $\varrho_{\tau}$ is the representation in antisymmetric tensors of degree $j$, we shall use the notation $d_{j}(\Theta)$ for $d^{2}(\Theta)$.

Since the restriction of $\Gamma_{|\tau|}$ to $\pi_{|\tau|}\left(p_{\tilde{\tau}}\right) \mathcal{R}^{\otimes|\tau|}$ is a multiple of $\varrho_{\tau}$, we have

$$
\begin{align*}
d^{\tau}(\Theta) & =\wedge\left\{m \in H_{i}^{\infty} \mid m H^{2}\left(\pi_{|\tau|}\left(p_{\bar{\tau}}\right) \Omega^{\otimes|\tau|}\right) \subset\right.  \tag{2.6}\\
& \left.\subset\left(\Gamma_{|\tau|}(\Theta) H^{2}\left(\pi_{|\tau|}\left(p_{\tilde{z}}\right) \Omega^{\otimes|\tau|}\right)\right)^{-}\right\} .
\end{align*}
$$

For the next lemma let $\tau^{\prime}, \tau$ be signatures with $\left|\tau^{\prime}\right|=n,|\tau|=n+1$ ( $n$ finite), $\ell\left(\tau^{\prime}\right)=\ell(\tau)=\operatorname{dim} \mathfrak{N}$ and such that denoting by $\tilde{\tau}^{\prime}$ and $\tilde{\tau}$ the signatures of length $n$, $n+1$, essentially equivalent to $\tau^{\prime}, \tau$, respectively, we have

Lemma 2.3. For $\Theta$ in $H^{\infty}(\mathscr{L}(\Omega))$ and $\tau^{\prime}, \tau$ as above, $d^{\tau}(\Theta)$ divides $d^{\tau}(\Theta)$.
Proof. Consider $m \in H_{i}^{\infty}$ such that

$$
m H^{2}\left(\pi_{n+1}\left(p_{\bar{z}}\right) \mathfrak{\Re}^{\otimes(n+1)}\right) \subset\left(\Gamma_{n+1}(\Theta) H^{2}\left(\pi_{n+1}\left(p_{\bar{z}}\right) \mathcal{S}^{\otimes(n+1)}\right)\right)^{-}
$$

It follows from Lemma 2.2. that

$$
\begin{gathered}
m\left(\bigvee_{k \in \Omega}\left(\left(\pi_{n}\left(p_{\tilde{z}^{\prime}}\right) T_{k}^{*} \pi_{n+1}\left(p_{\mathfrak{z}}\right)\right) \otimes I_{H^{2}}\right) H^{2}\left(\Omega^{n+1}\right)\right)^{-} \subset \\
\subset\left(\left(\pi_{n}\left(p_{\tilde{z}^{\prime}}\right) \otimes I_{H^{2}}\right) \Gamma_{n}(\Theta) H^{2}\left(\boldsymbol{\Omega}^{\otimes n}\right)\right)^{-}=\left(\Gamma_{n}(\Theta) H^{2}\left(\pi_{n}\left(p_{\mathfrak{r}^{\prime}}\right) \Omega^{\otimes n}\right)\right)^{-}
\end{gathered}
$$

and hence by Lemma 2.1

$$
m H^{2}\left(\pi_{n}\left(p_{\hat{\tau}^{\prime}}\right) \mathcal{R}^{\otimes n}\right) \subset\left(\Gamma_{n}(\Theta) H^{2}\left(\pi_{n}\left(p_{\tilde{\mathrm{r}}^{\prime}}\right) \mathcal{\Omega}^{\otimes n}\right)\right)^{-}
$$

so that by (2.6) $d^{\tau^{\prime}}$ divides $m$.
Q.E.D.

Let us also record the following simple fact for further use.
Remark 2.4. Let $\mathfrak{X}_{i}, \mathfrak{V}_{i}(i=1,2)$ be Hilbert spaces, $A_{i} \in H^{\infty}\left(\mathscr{L}\left(\mathfrak{X}_{i}, \mathfrak{Y}_{i}\right)\right)$, $B \in H^{\infty}\left(\mathscr{L}\left(\mathfrak{X}_{1} \otimes \mathfrak{X}_{2}, \mathfrak{Y}_{1} \otimes \mathfrak{Y}_{2}\right)\right), B(z)=A_{1}(z) \otimes A_{2}(z)(z \in D)$ and suppose $f_{i} \in\left(A_{i} H^{2}\left(\mathfrak{X}_{i}\right)\right)^{-} \cap$ $\cap H^{\infty}\left(\mathfrak{Y}_{i}\right)$. Then we have $f_{1} \otimes f_{2} \in\left(B H^{2}\left(\mathfrak{X}_{1} \otimes \mathfrak{X}_{2}\right)\right)^{-}$. Indeed, consider $h_{i}^{(n)} \in H^{\infty}\left(\mathfrak{X}_{i}\right)$ such that

$$
\lim _{n \rightarrow \infty}\left\|A_{i} h_{i}^{(n)}-f_{i}\right\|=0 \quad \text { in } \quad H^{2}\left(\mathfrak{Y}_{i}\right)
$$

Then in

$$
H^{2}\left(\mathfrak{Y}_{1} \otimes \mathfrak{Y}_{2}\right)
$$

we have

$$
\lim _{n \rightarrow \infty}\left\|B\left(h_{1}^{(n)} \otimes h_{2}^{(m)}\right)-f_{1} \otimes A_{2} h_{2}^{(m)}\right\|=0
$$

and

$$
\lim _{m \rightarrow \infty}\left\|f_{1} \otimes f_{2}-f_{1} \otimes A_{2} h_{2}^{(m)}\right\|=0
$$

which is the desired result.

For the following theorem consider $\Theta \in H^{\infty}(\mathscr{L}(\Omega)), \quad \Theta^{\prime} \in H^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}\right)\right)$ and suppose there are $A \in H^{\infty}\left(\mathscr{L}\left(\Omega, \Omega^{\prime}\right)\right), B \in H^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}, \Omega\right)\right)$ such that the following set of relations holds

$$
\left\{\begin{array}{l}
A \Theta H^{2}(\Omega) \subset\left(\Theta^{\prime} H^{2}\left(\Omega^{\prime}\right)\right)^{-}  \tag{2.7}\\
B \Theta^{\prime} H^{2}\left(\Omega^{\prime}\right) \subset\left(\Theta H^{2}(\Omega)\right)^{-} \\
B A H^{2}(\Omega) \vee \Theta H^{2}(\Omega)=H^{2}(\Omega)
\end{array}\right.
$$

Theorem 2.5. Let $\Theta, \Theta^{\prime}, A, B$ be as before and suppose (2.7) holds. Let further $\tau, \tau^{\prime}$ be essentially equivalent signatures with $l(\tau)=\operatorname{dim} \Omega, l\left(\tau^{\prime}\right)=\operatorname{dim} \Omega^{\prime},|\tau|<\infty$ and ${ }_{\varepsilon}(\tau)=』\left(\tau^{\prime}\right) \leqq \min \left(\operatorname{dim} \boldsymbol{\Omega}, \operatorname{dim} \Omega^{\prime}\right)$. Then $d^{\tau}(\Theta)$ divides $d^{\tau^{\prime}}\left(\Theta^{\prime}\right)$.

Proof. If $d^{t^{\prime}}\left(\Theta^{\prime}\right)=0$, the assertion of the theorem is obvious, so assume $d^{r^{\prime}}\left(\Theta^{\prime}\right)=m \in H_{i}^{\infty}$. Let $\tilde{\tau}$ denote the signature of length $n=|\tau|$ that is essentially equivalent to $\tau$.

Consider $f_{1}, f_{2}, \ldots, f_{n} \in H^{\infty}(\Omega), g_{1}, g_{2}, \ldots, g_{n} \in H^{\infty}(\Omega)$ and

$$
\begin{equation*}
s=\left(\pi_{n}\left(p_{\tilde{i}}\right) \otimes I_{H^{2}}\right)\left(\left(B A f_{1}+\Theta g_{1}\right) \otimes \ldots \otimes\left(B A f_{n}+\Theta g_{n}\right)\right) \tag{2.8}
\end{equation*}
$$

Using (2.7) it is easily seen that the elements $s$ form a total subset of $H^{2}\left(\pi_{n}\left(p_{\tilde{z}}\right) \mathfrak{\Omega}^{\otimes n}\right)$, so that it will be sufficient to prove that

$$
\begin{equation*}
m s \in\left(\Gamma_{n}(\Theta) H^{2}\left(\pi_{n}\left(p_{\bar{\tau}}\right) \Omega^{\otimes n}\right)\right)^{-} \tag{2.9}
\end{equation*}
$$

Now, $s$ is a finite sum of elements of the form

$$
\begin{equation*}
r=\left(\left(\pi_{n}\left(p_{\tau}\right) \pi_{n}(\sigma)\right) \otimes I_{H^{2}}\right)\left(B A f_{1}^{\prime} \otimes \ldots \otimes B A f_{j}^{\prime} \otimes \Theta g_{1}^{\prime} \otimes \ldots \otimes \Theta g_{n-\jmath}^{\prime}\right) \tag{2.10}
\end{equation*}
$$

where $0 \leqq j \leqq n, \sigma \in \Im_{n}$ and $f_{i}^{\prime}, g_{i}^{\prime}$ are some of the $f^{\prime}$ and $g$. Thus to prove (2.9) it will be enough to show that

$$
\begin{equation*}
m r \in\left(\Gamma_{n}(\Theta) H^{2}\left(\pi_{n}\left(p_{\bar{\tau}}\right) \Omega^{\otimes n}\right)\right)^{-} \tag{2.11}
\end{equation*}
$$

Because $\sum_{\gamma \in \hat{\epsilon}_{j}} p_{\gamma}=1$ and $\Theta_{j}$ is considered as a subgroup of $\Theta_{n}(j \leqq n)$, we have $\sum_{\gamma \in \tilde{E}_{j}} p_{\tilde{\imath}} p_{\gamma}=p_{\tilde{\tau}}$ and $p_{\tilde{\tau}} p_{\gamma} \neq 0$ if and only if the restriction of the representation of signature $\tilde{\tau}$ to $\mathbb{S}_{j}$ contains the representation of signature $\gamma$. So, $p_{\tilde{\tau}} p_{\gamma} \neq 0$ if and only if there are $\gamma_{k} \in \hat{ভ}_{k}(j<k<n)$ such that

$$
\begin{equation*}
\gamma \prec \gamma_{j+1} \prec \ldots \prec \gamma_{n-1} \prec \tilde{\tau} . \tag{2.12}
\end{equation*}
$$

Hence denoting by $\check{\gamma}$ the signature of length $\operatorname{dim} \boldsymbol{\Omega}^{\prime}$ that is essentially equivalent to $\gamma$, using Lemma 2.3 several times we conclude that $d^{\breve{\gamma}}\left(\Theta^{\prime}\right)$ divides $d^{\tau}\left(\Theta^{\prime}\right)=m$.

Now we have:

$$
\begin{gathered}
m r= \\
=\left(\pi_{n}\left(p_{\tilde{i}}\right) \pi_{n}(\sigma) \otimes I_{H^{2}}\right) \sum_{\substack{\gamma \in \hat{E}_{j \neq 0} \\
P_{i}^{\prime} P_{\gamma} \neq 0}}\left(m\left(\pi_{j}\left(p_{\gamma}\right) \otimes I_{H^{2}}\right)\left(B A f_{1}^{\prime} \otimes \ldots \otimes B A f_{j}^{\prime}\right)\right) \otimes\left(\Theta g_{1}^{\prime} \otimes \ldots \otimes \Theta g_{n-j}^{\prime}\right) .
\end{gathered}
$$

To end the proof it will be sufficient to show that

$$
m\left(\pi_{j}\left(p_{\gamma}\right) \otimes I_{H^{2}}\right)\left(B A f_{1}^{\prime} \otimes \ldots \otimes B A f_{j}^{\prime}\right) \quad \text { is in } \quad\left(\Gamma_{j}(\Theta) H^{2}\left(\Omega^{\otimes j}\right)\right)^{-}
$$

because then using Remark 2.4 we will have that $m r$ is in

$$
\left(\left(\pi_{n}\left(p_{\overline{\mathfrak{z}}}\right) \pi_{n}(\sigma) \otimes I_{H^{2}}\right) \Gamma_{n}(\Theta) H^{2}\left(\Omega^{\otimes n}\right)\right)^{-}=\left(\Gamma_{n}(\Theta) H^{2}\left(\pi_{n}\left(p_{\bar{i}}\right) \Omega^{\otimes n}\right)\right)^{-}
$$

which is the desired result.
Now further $m\left(\pi_{j}\left(p_{\gamma}\right) \otimes I_{H^{2}}\right)\left(A f_{1}^{\prime} \otimes \ldots \otimes A f_{j}^{\prime}\right)$ is in $d^{\bar{y}}\left(\Theta^{\prime}\right) H^{2}\left(\pi_{j}\left(p_{\gamma}\right) \Omega^{\prime \otimes j}\right)$, since $d^{\check{y}}\left(\Theta^{\prime}\right)$ divides $m$, and hence is in $\left(\Gamma_{j}\left(\Theta^{\prime}\right) H^{2}\left(\pi_{j}\left(p_{y}\right) \Omega^{\prime \otimes j}\right)\right)^{-} \subset\left(\Gamma_{j}\left(\Theta^{\prime}\right) H^{2}\left(\boldsymbol{\Omega}^{\prime \otimes j}\right)\right)^{-}$. Thus it will be sufficient to prove that

$$
\left(\Gamma_{j}(B) \Gamma_{j}\left(\Theta^{\prime}\right) H^{2}\left(\Omega^{\prime \otimes j}\right)\right)^{-} \subset\left(\Gamma_{j}(\Theta) H^{2}\left(\Omega^{\otimes j}\right)\right)^{-}
$$

in order that

$$
m\left(\pi_{j}\left(p_{\gamma}\right) \otimes I_{H^{2}}\right)\left(B A f_{1}^{\prime} \otimes \ldots \otimes B A f_{j}^{\prime}\right) \in\left(\Gamma_{j}(\Theta) H^{2}\left(\Omega^{\otimes j}\right)\right)^{-}
$$

To this end remark that the elements of the form $B \Theta^{\prime} h_{1} \otimes \ldots \otimes B \Theta^{\prime} h_{j}$ with $h_{i} \in H^{\infty}\left(\boldsymbol{\Omega}^{\prime}\right)$ are total in $\left(\Gamma_{j}(B) \Gamma_{j}\left(\Theta^{\prime}\right) H^{2}\left(\Omega^{\prime \otimes j}\right)\right)^{-}$and

$$
B \Theta^{\prime} h_{1} \otimes \ldots \otimes B \Theta^{\prime} h_{j} \in\left(\Gamma_{j}(\Theta) H^{2}\left(\Omega^{\otimes j}\right)\right)^{-}
$$

because fo (2.7) and Remark 2.4.
Q.E.D.

## § 3. Applications to quasi-similar $C_{0}$ operators

The following Proposition is an easy application of Theorem 2.5.
Proposition 3.1. Let $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$, $\Theta^{\prime} \in H_{i}^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}\right)\right)$ and let $\tau, \tau^{\prime}$ be essentially equivalent signatures with $l(\tau)=\operatorname{dim} \Omega, \quad l\left(\tau^{\prime}\right)=\operatorname{dim} \mathcal{R}^{\prime}$ and $\_(\tau)=\_\left(\tau^{\prime}\right) \leqq$ $\leqq \min \left(\operatorname{dim} \Omega, \operatorname{dim} \Omega^{\prime}\right)$. If $S(\Theta)$ and $S\left(\Theta^{\prime}\right)$ are quasi-similar, we have

$$
\begin{equation*}
d^{\tau}(\Theta)=d^{\tau^{\prime}}\left(\Theta^{\prime}\right) \tag{3.1}
\end{equation*}
$$

Proof. Let $X$ and $Y$ be two quasi-affinities such that $S\left(\Theta^{\prime}\right) X=X S(\Theta)$ and $S(\Theta) Y=Y S\left(\Theta^{\prime}\right)$. From the lifting theorem (see (1.8-11)) it follows that we can find $A \in H^{\infty}\left(\mathscr{L}\left(\Omega, \Omega^{\prime}\right)\right)$ and $B \in H^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}, \Omega\right)\right)$ such that

$$
\begin{equation*}
X=P_{5\left(\theta^{\prime}\right)} A\left|\mathfrak{H}(\Theta), \quad Y=P_{\mathfrak{5}(\boldsymbol{\theta})} B\right| \mathfrak{G}\left(\Theta^{\prime}\right) \tag{3.2}
\end{equation*}
$$

$$
\begin{equation*}
A \Theta H^{2}(\Omega) \subset \Theta^{\prime} H^{2}\left(\Omega^{\prime}\right), \quad B \Theta^{\prime} H^{2}\left(\Omega^{\prime}\right) \subset \Theta H^{2}(\Omega) \tag{3.3}
\end{equation*}
$$

and

$$
\begin{equation*}
A B H^{2}\left(\Omega^{\prime}\right) \vee \Theta^{\prime} H^{2}\left(\Omega^{\prime}\right)=H^{2}\left(\Omega^{\prime}\right), \quad B A H^{2}(\Omega) \vee \Theta H^{2}(\Omega)=H^{2}(\Omega) \tag{3.4}
\end{equation*}
$$

so that the assumptions of Theorem 2.5 are satisfied. It follows that $d^{\tau}(\Theta)$ divides $d^{\tau^{\prime}}\left(\Theta^{\prime}\right)$ and $d^{t^{t}}\left(\Theta^{\prime}\right)$ divides $d^{t}(\Theta)$ and this proves (3.1).
Q.E.D.

Let $T$ be any operator unitarily equivalent to some $S(\Theta)$ with a pure $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$. It is easy to see that the functions $d^{\tau}(\Theta)$ and $d_{j}(\Theta)$ depend only on $T$ and not on the particular function $\Theta$, so we shall denote them by $d^{\tau}(T)$ and $d_{j}(T)$, respectively.

Corollary 3.2. If $T$ and $T^{\prime}$ are two quasisimilar $C_{0}$ operators and $\mathfrak{D}_{T}=\mathfrak{D}_{T^{\prime}}$, then $d^{\tau}(T)=d^{\tau}\left(T^{\prime}\right)$ for each $\tau$ with $l(\tau)=D_{T}$.

Proof. $T$ and $T^{\prime}$ are unitarily equivalent to $S(\Theta)$ and $S\left(\Theta^{\prime}\right)$, respectively, where $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega)), \Theta^{\prime} \in H_{i}^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}\right)\right)$ with $\operatorname{dim} \Omega=\operatorname{dim} \Omega^{\prime}=\mathfrak{D}_{T}$. The corollary obviously follows from Proposition 3.1.
Q.E.D.

Consider now a $C_{0}$ operator $T$ with Jordan model $S=S\left(m_{1}\right) \oplus S\left(m_{2}\right) \oplus \ldots$. If $\mathfrak{D}_{S}<\mathfrak{D}_{T}$ we shall put $m_{j}=1$ for $\mathfrak{D}_{S}<j \leqq \mathfrak{D}_{T}$. So we have

$$
\begin{equation*}
S=\oplus_{j=1}^{D_{T}} S\left(m_{j}\right) \tag{3.5}
\end{equation*}
$$

Corollary 3.3. For any $C_{0}$ operator $T$ and any signature $\tau=\left(t_{1} \geqq t_{2} \geqq \ldots\right)$, $|\tau|<\infty, l(\tau)=\delta_{T}$, we have

$$
\begin{equation*}
d^{\tau}(T)=m_{1}^{t_{1}} m_{2}^{t_{2}}, \ldots, m_{n}^{t_{n}}, \quad n=s(\tau) \tag{3.6}
\end{equation*}
$$

Proof. We have only to apply Proposition 3.1 to $\Theta$ coinciding with the characteristic function of $T$ and to

$$
\Theta^{\prime}=\operatorname{diag}\left(m_{1}, m_{2}, \ldots\right) \in H_{i}^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}\right)\right) \quad \text { with } \quad \operatorname{dim} \Omega^{\prime}=\mathfrak{b}_{T}
$$

Since $\tau=\left(t_{1} \geqq t_{2} \geqq \ldots\right)$ represents the highest weight to the representation $\varrho_{\tau}$ (see [18], [6] to the finite-dimensional and [1] for the infinite-dimensional case) it is immediate that:

$$
d^{\tau}\left(\Theta^{\prime}\right)=m_{1}^{t_{1}}, \ldots, m_{n}^{t_{n}}
$$

Q.E.D.

Corollary 3.4. For any $C_{0}$ operator $T$, the functions $m_{j}$ appearing in the Jordan model can be computed as

$$
\begin{equation*}
m_{j}=d_{j}(T) / d_{j-1}(T), \quad 1 \leqq j \leqq \mathfrak{D}_{T} \quad \text { where } \quad d_{0}(T)=1 \tag{3.7}
\end{equation*}
$$

Proof. The preceding Corollary gives for $\tau_{j}=(1, \ldots, 1,0, \ldots)$ (with $j$ nonzero terms)

$$
d_{j}(T)=d^{\tau j}(T)=m_{1} \ldots m_{j}, \quad j \leqq \mathfrak{D}_{T}
$$

so relation (3.7) becomes obvious.
Q.E.D.

Since the quasisimilarity class of a $C_{0}$ operator is determined by the Jordan model, Corollary 3.4 shows that a $C_{0}$ operator $T$ is determined up to quasisimilarity by the least inner multiples of the exterior powers of any function coinciding with the characteristic function of $T$. This enables us to prove the following theorem.

Theorem 3.5. Let $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega)), \Theta^{\prime} \in H_{i}^{\infty}\left(\mathscr{L}\left(\Omega^{\prime}\right)\right)$ be such that $d_{1}(\Theta) \neq 0$ and $\operatorname{dim} \Omega=\operatorname{dim} \Omega^{\prime}$. If $S(\Theta)$ and $S\left(\Theta^{\prime}\right)$ are quasisimilar then $S\left(\varrho_{\tau}(\Theta)\right)$ and $S\left(\varrho_{\tau}\left(\Theta^{\prime}\right)\right)$ are quasisimilar for each signature $\tau$ such that $l(\tau)=\operatorname{dim} \Omega,|\tau|<\infty$.

Proof. By Corollary 3.4 we have only to show that $d_{j}\left(\varrho_{\tau}(\Theta)\right)=d_{j}\left(\varrho_{\tau}\left(\Theta^{\prime}\right)\right)$ for each $j \leqq \operatorname{dim} \mathfrak{\Omega}^{\tau}$. Let $\tau_{j}=\left(1,1, \ldots, 1,0, \ldots\right.$ ) (with $j$ nonzero terms), $\ell\left(\tau_{j}\right)=\operatorname{dim} \mathcal{\Omega}^{\boldsymbol{T}}$.

The representation $\varrho_{\tau} \circ \varrho_{\tau}$ of $\mathscr{U}(\Omega)$ is a subrepresentation of the representation of $\mathscr{U}(\Omega)$ on $\Omega^{\otimes j|\tau|}$ and hence a finite direct sum of representations $\varrho_{\tau^{\prime}}$, with $\ell\left(\tau^{\prime}\right)=$ $=\operatorname{dim} \Omega,\left|\tau^{\prime}\right|<\infty$ :

$$
\begin{equation*}
\varrho_{\tau_{j}} \circ \varrho_{\tau}=\underset{\tau^{\prime}}{\oplus} \varrho_{\tau^{\prime}} \tag{3.8}
\end{equation*}
$$

From (3.8) it follows then that

$$
\varrho_{\tau_{j}}\left(\varrho_{\tau}(\Theta)\right)=\underset{\tau^{\prime}}{\oplus} \varrho_{\tau^{\prime}}(\Theta), \quad \varrho_{\tau_{j}}\left(\varrho_{\tau}\left(\Theta^{\prime}\right)\right)=\underset{\mathfrak{\tau}^{\prime}}{\oplus} \varrho_{\tau^{\prime}}\left(\Theta^{\prime}\right)
$$

and hence $d_{j}\left(\varrho_{\tau}(\Theta)\right)$ is the least inner common multiple of the $d^{r^{\prime}}(\Theta)$ and $d_{j}\left(\varrho_{\tau}\left(\Theta^{\prime}\right)\right)$ the least inner common multiple of the $d^{t^{\prime}}\left(\Theta^{\prime}\right)$. Since $d^{t^{\prime}}(\Theta)=d^{r^{\prime}}\left(\Theta^{\prime}\right)$ by Proposition 3.1, we infer that $d_{j}\left(\varrho_{\tau}(\Theta)\right)=d_{j}\left(\varrho_{\tau}\left(\Theta^{\prime}\right)\right)$.


#### Abstract

Q.E.D.


## $\S$ 4. Defect operators of $C_{0}$ contractions

For an operator $A \in \mathscr{L}(\mathfrak{R})$ and a closed subspace $\mathfrak{P} \subset \mathfrak{\Omega}$ we consider

$$
\gamma[A, \mathfrak{M}]=\inf _{\substack{k \in \mathfrak{M} \\\|k\|=1}}\|A k\|, \quad \gamma_{j}(A)=\sup _{\operatorname{codim} \mathfrak{N}=j-1} \gamma[A, \mathfrak{M}] .
$$

As is known from the minimax principle, $\gamma_{j}(A)(1 \leqq j \leqq \operatorname{dim} \Omega)$ are eigenvalues of $\left(A^{*} A\right)^{1 / 2}$ in increasing order. In case $\operatorname{dim} \Omega<\infty$ all eigenvalues of $\left(A^{*} A\right)^{1 / 2}$ repeated according to their multiplicity appear in the sequence of the $\gamma_{j}(A)$. In case $\operatorname{dim} \Omega=\infty$, $\gamma_{1}(A)$ is the least eigenvalue of $\left(A^{*} A\right)^{1 / 2}$, discrete eigenvalues smaller than the least essential eigenvalue appear in increasing order repeated according to their multiplicity and the sequence becomes stationary if the least essential eigenvalue of $\left(A^{*} A\right)^{1 / 2}$ is reached.

For the next two lemmas, $\tau_{j}$ denotes the signature

$$
\tau_{j}=(1, \ldots, 1,0 \ldots), \quad l\left(\tau_{j}\right)=\operatorname{dim} \Omega, \quad s\left(\tau_{j}\right)=j
$$

Lemma 4.1. Let $A \in \mathscr{L}(\Omega)$ and $\tau_{j}$ be as above. Then we have:

$$
\begin{equation*}
\gamma_{1}\left(\varrho_{\tau_{j}}(A)\right)=\gamma_{1}(A) \gamma_{2}(A) \ldots \gamma_{j}(A) \tag{4.1}
\end{equation*}
$$

Proof. Remark first that applying $\varrho_{\tau_{j}}$, to the polar decomposition of $A$ we get the polar decomposition of $\varrho_{\tau_{j}}(A)$, so we can suppose $A$ is positive. Moreover, in
view of the minimax definition of $\gamma_{j}$, we have $\left|\gamma_{j}(A)-\gamma_{j}(B)\right| \leqq\|A-B\|$, and thus by continuity it will be sufficient to consider the case when $A \geqq 0$ has finite spectrum.

In this case, $\varrho_{\tau_{j}}$ being the representation in antisymmetric tensors of degree $j, \varrho_{\tau_{j}}(A)$ has finite spectrum, the eigenvalues being products $\lambda_{1} \ldots \lambda_{j}$ of eigenvalues of $A$, a given eigenvalue appearing in such a product at most a number of times equal to its multiplicity. Clearly $\gamma_{1}(A) \ldots \gamma_{j}(A)$ is then the least eigenvalue of $\varrho_{\tau_{j}}(A)$.
Q.E.D.

Lemma 4.2. Let $T$ be a $C_{0}$ operator, let $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$ coincide with the characteristic function of $T$ and let $\left\{m_{j}\right\}_{j=1}^{b_{T}}$ be inner functions for the Jordan model of $T$ with $m_{j} \equiv 1$ for $\mu_{T}<j \leqq D_{T}$. Then we have

$$
\begin{equation*}
\gamma_{1}(\Theta(\lambda)) \ldots \gamma_{j}(\Theta(\lambda)) \geqq\left|m_{1}(\lambda) \ldots m_{j}(\lambda)\right| \tag{4.2}
\end{equation*}
$$

where $1 \leqq j \leqq D_{T}$ and $\lambda \in D$.
Proof. In view of Corollary 3.3, $m_{1} \ldots m_{j}$ is the least inner multiple of $\varrho_{\tau_{j}}(\Theta) \in$ $\in H_{i}^{\infty}\left(\mathscr{L}\left(\Omega^{\boldsymbol{\tau}}\right)\right)$. Hence there is a contractive function $\Omega \in H^{\infty}\left(\mathscr{L}\left(\boldsymbol{\Omega}^{\tau_{j}}\right)\right)$ such that

$$
\Omega(\lambda) \varrho_{\tau_{j}}(\Theta(\lambda))=m_{1}(\lambda) \ldots m_{j}(\lambda) I_{\Omega} \tau_{j}
$$

Since $\|\Omega(\lambda)\| \leqq 1$ this clearly implies

$$
\gamma_{1}\left(\varrho_{\tau_{j}}(\Theta(\lambda))\right) \geqq\left|m_{1}(\lambda) \ldots m_{j}(\lambda)\right|
$$

and by Lemma 4.1

$$
\gamma_{1}\left(\varrho_{t_{j}}(\Theta(\lambda))\right) \leqq \gamma_{1}(\Theta(\lambda)) \ldots \gamma_{j}(\Theta(\lambda)),
$$

which gives the desired inequality.
Q.E.D.

Proposition 4.3. Let $T$ be a $C_{0}$ operator acting on $\mathfrak{G}$ and $\left\{m_{j}\right\}_{j=1}^{\infty}$ inner functions for the Jordan model of $T$, with $m_{j} \equiv 1$ in case $\mu_{T}<j$.
a) If $\sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|\right)<\infty$, then $\operatorname{tr}\left(I-T^{*} T\right)<\infty$.
b) If $\lim _{j \rightarrow \infty}\left|m_{j}(0)\right|=1$, then $I-T^{*} T$ is compact.

Proof. a) The assumptions are that the Jordan model $S=S\left(m_{1}\right) \oplus S\left(m_{2}\right) \oplus \ldots$ is a weak contraction ([13] ch. VIII) since $\operatorname{tr}\left(I-S^{*} S\right)=\sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|^{2}\right) \leqq$ $\leqq 2 \sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|\right)<\infty$. As usual for weak contractions there will be no loss of generality to assume that $m_{j}(0) \neq 0$ (one uses a conformal automorphism of the unit disc as in [13] ch. VIII). Thus the infinite product $\prod_{j=1}^{\infty}\left|m_{j}(0)\right|$ converges to some $c>0$. Hence by Lemma 4.2 for $\Theta$ the characteristic function of $T$, we infer that

$$
\prod_{1 \leqq j \leqq \mathrm{o}_{\boldsymbol{T}}} \gamma_{j}(\Theta(0))>0 .
$$

Since in case $\boldsymbol{D}_{T}=\infty$ this implies $\lim _{j \rightarrow \infty} \gamma_{j}(\Theta(0))=1$, it follows that

$$
\operatorname{tr}\left(I_{\mathfrak{D}_{T}}-\Theta(0)^{*} \Theta(0)\right)=\sum_{1 \equiv j \leqslant \mathbf{D}_{\boldsymbol{T}}}\left(1-\gamma_{j}(\Theta(0))^{2}\right)
$$

and

$$
\sum_{1 \leqq j \leqq b_{r}}\left(1-\gamma_{j}(\Theta(0))^{2}\right)<\infty
$$

since

$$
\prod_{\mathbf{1} \leq j \leq \mathrm{D}_{T}} \gamma_{j}(\Theta(0))>0 . \quad \text { But } \quad I_{\mathcal{D}_{T}}-\Theta(0)^{*} \Theta(0)=D_{T}^{2} \mid \mathfrak{D}_{T}
$$

so that $\operatorname{tr}\left(I-T^{*} T\right)<\infty$.
b) The proof is quite similar to that of a), so we can be brief in details. Again we may suppose $T$ is invertible and hence $m_{j}(0) \neq 0$. Then $\lim _{j \rightarrow \infty}\left|m_{j}(0)\right|=1$ gives

$$
\lim _{j \rightarrow \infty}\left|m_{1}(0) \ldots m_{j}(0)\right|^{1 / j}=1
$$

Using Lemma 4.2 this implies

$$
\lim _{j \rightarrow \infty}\left(\gamma_{1}(\Theta(0)) \ldots \gamma_{j}(\Theta(0))\right)^{1 / j}=1
$$

so that $\lim _{j \rightarrow \infty} \gamma_{j}(\Theta(0))=1$ which gives that $I-T^{*} T$ is compact.

Remark 4.4. As we shall see in $\S 8$ the converse of 4.3 a) is also true. For 4.3 b) the converse is in general false. An example can be constructed as follows.

Let $\mu$ be a finite non-negative measure on $[0,2 \pi]$, singular with respect to Lebesgue measure and without atoms. Consider the inner functions

$$
m_{j, n}(\lambda)=\exp \left[-\int_{2 \pi(j-1) / n}^{2 \pi j / n} \frac{e^{i t}+\lambda}{e^{i t}-\lambda} d \mu(t)\right], \quad 1 \leqq j \leqq n
$$

and the operators

$$
T=\oplus_{n=1}^{\infty}\left(\oplus_{j=1}^{n} S\left(m_{j, n}\right)\right), \quad S=S\left(m_{1,1}\right) \oplus S\left(m_{1,1}\right) \oplus \ldots
$$

Then $S$ is the Jordan model of $T, I-T^{*} T$ is compact and $\left|m_{1,1}(0)\right|,\left|m_{1,1}(0)\right|, \ldots$ tends to $\left|m_{1,1}(0)\right|<1$.

Proposition 4.5. Let $T$ be a $C_{0}$ operator, let $\left\{m_{j}\right\}_{j=1}^{\infty}$ be inner functions for the Jordan model of $T\left(m_{j} \equiv 1\right.$ in case $\left.\mu_{T}<j\right)$ and let $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$ coincide with the characteristic function of $T$. Suppose moreover $m_{1}(0) \neq 0$ and $n \in \mathbf{N}$ is such that $\left|m_{n}(0)\right|<$ $<\lim _{j \rightarrow \infty}\left|m_{j}(0)\right|$. Then the following conditions are equivalent:
(i) $\left|m_{1}(0) \ldots m_{n}(0)\right|=\gamma_{1}(\Theta(0)) \ldots \gamma_{n}(\Theta(0))$,
(ii) $T$ is unitarily equivalent to $T_{1} \oplus T_{2}$, where $\mathfrak{b}_{r_{1}}=n$ and $T_{1}, T_{2}$ are quasisimilar to $S\left(m_{1}\right) \oplus \ldots \oplus S\left(m_{n}\right)$ and respectively to $S\left(m_{n+1}\right) \oplus S\left(m_{n+2}\right) \oplus \ldots$.

Proof. (i) $\Rightarrow$ (ii). The condition $0 \neq\left|m_{n}(0)\right|<\lim _{j \rightarrow \infty}\left|m_{j}(0)\right|$ implies that $\gamma_{n}(\Theta(0))$ is less than the least essential eigenvalue of $\left(\Theta(0)^{*} \Theta(0)\right)^{1 / 2}$, for otherwise we would have $\gamma_{n}(\Theta(0))=\gamma_{n+1}(\Theta(0))=\ldots$ which in view of Lemma 4.2 would imply $\lim _{j \rightarrow \infty}\left|m_{j}(0)\right| \leqq$ $\leqq \gamma_{n}(\Theta(0))$ and hence $\left|m_{n}(0)\right|<\gamma_{n}(\Theta(0))$ which when combined with (i) would give $\left|m_{1}(0) \ldots m_{n-1}(0)\right|>\gamma_{1}(\Theta(0)) \ldots \gamma_{n-1}(\Theta(0))$, contradicting Lemma 4.2. Thus replacing $\Theta$ by some equivalent inner operator-valued function in $H_{i}^{\infty}(\mathscr{L}(\Omega))$ we may assume there is an orthonormal set $\left\{e_{1}, \ldots, e_{n}\right\}$ in $\Omega$ such that $\Theta(0) e_{j}=\gamma_{j}(\Theta(0)) e_{j}$ for $1 \leqq j \leqq n$. Consider $f=\pi_{n}\left(p_{\tau_{n}}\right)\left(e_{1} \oplus \ldots \oplus e_{n}\right)$. Then

$$
\varrho_{\tau_{n}}(\Theta(0)) f=\gamma_{1}(\Theta(0)) \ldots \gamma_{n}(\Theta(0)) f
$$

and since $p_{\tau_{n}}=(n!)^{-1} \sum_{\sigma \in \Xi_{n}} \varepsilon(\sigma) \sigma(\varepsilon(\sigma)$ is the sign of the permutation $\sigma)$, we have $f \neq 0$. But $\Omega \varrho_{\tau_{n}}(\Theta)=m_{1} \ldots m_{n} I_{\Omega_{n}}$ for some contractive $\Omega$, and we infer $\|\Omega(0) f\|=\|f\|$ so that $\Omega(\lambda) f=\mu f$ for some constant $\mu,|\mu|=1$. This in turn implies $\varrho_{\tau_{n}}(\Theta(\lambda)) f=$ $=\mu^{-1} m_{1}(\lambda) \ldots m_{n}(\lambda) f$ for all $\lambda \in D$. In view of the known properties of $p_{\tau_{n}}$ this last equality implies that $\mathfrak{B}=\mathbf{C} e_{1}+\ldots+\mathbf{C} e_{n}$ is invariant for $\Theta(\lambda)$ for all $\lambda \in D$. Since $\Theta$ is two-sided inner we infer that $\mathfrak{B}$ is a reducing subspace for $\Theta(\lambda), \lambda \in D$. Hence $\Theta=\Theta_{1} \oplus \Theta_{2}$ where $\Theta_{1}=\Theta\left|\mathfrak{B}, \Theta_{2}=\Theta\right| \Omega \ominus \mathfrak{B}$.

Thus we define $T_{i}=S\left(\Theta_{i}\right)$ for $i=1,2$ and clearly $T$ is unitarily equivalent to $T_{1} \oplus T_{2}$ and $\mathrm{D}_{T_{1}}=n$. Remark also that $\varrho_{\tau_{n}}\left(\Theta_{1}\right)$ coincides with $m_{1} \ldots m_{n}$. Let $S\left(m_{1}^{\prime}\right) \oplus \ldots$ $\ldots \oplus S\left(m_{n}^{\prime}\right)$ and $S\left(m_{1}^{\prime \prime}\right) \oplus S\left(m_{2}^{\prime \prime}\right) \oplus \ldots$ be the Jordan models of $T_{1}$ and $T_{2}$ (we do not exclude the possibility that some $m_{j}^{\prime}$ or $m_{j}^{\prime \prime}$ be 1 ). Then we have:

$$
\begin{equation*}
m_{1} \ldots m_{n}=m_{1}^{\prime} \ldots m_{n}^{\prime}=\bigvee_{k=0}^{n} m_{1}^{\prime} \ldots m_{k}^{\prime} m_{1}^{\prime \prime} \ldots m_{n-k}^{\prime \prime} \tag{4.3}
\end{equation*}
$$

(use for instance Proposition 3.1 with $\tau=\tau_{n}$ ). From 4.3 we infer that $m_{1}^{\prime} \ldots m_{n-1}^{\prime} m_{1}^{\prime \prime}$ divides $m_{1}^{\prime} \ldots m_{n}^{\prime}$ and hence $m_{1}^{\prime \prime}$ divides $m_{n}^{\prime}$. Thus $S\left(m_{1}^{\prime}\right) \oplus \ldots \oplus S\left(m_{n}^{\prime}\right) \oplus S\left(m_{1}^{\prime \prime}\right) \oplus$ $\oplus S\left(m_{2}^{\prime \prime}\right) \oplus \ldots$ is the Jordan model of $T_{1} \oplus T_{2}$ and hence $m_{j}^{\prime}=m_{j}, m_{k}^{\prime \prime}=m_{n+k}$ ( $1 \leqq j \leqq n, k=1,2, \ldots$ ). This ends the proof of (i) $\Rightarrow$ (ii).
(ii) $\Rightarrow$ (i). Let $\Theta_{1}, \Theta_{2}$ coincide with the characteristic functions of $T_{1}, T_{2}$. Then $\varrho_{\tau_{n}}\left(\Theta_{1}\right)$ coincides with $m_{1} \ldots m_{n}$ so that $\gamma_{1}\left(\Theta_{1}(0)\right) \ldots \gamma_{n}\left(\Theta_{1}(0)\right)=\gamma_{1}\left(\tau_{n}\left(\Theta_{1}(0)\right)\right)=$ $=\left|m_{1}(0) \ldots m_{n}(0)\right|$ (use Proposition 3.1 for instance and then Lemma 4.1). Now clearly $\gamma_{j}\left(\Theta_{1}(0)\right) \geqq \gamma_{j}(\Theta(0))$ and hence $\gamma_{1}(\Theta(0)) \ldots \gamma_{n}(\Theta(0)) \leqq\left|m_{1}(0) \ldots m_{n}(0)\right|$ which in view of Lemma 4.2 gives $\gamma_{1}(\Theta(0)) \ldots \gamma_{n}(\Theta(0))=\left|m_{1}(0) \ldots m_{n}(0)\right|$.
Q.E.D

Remark 4.6. If $T$ is a contraction and $\Theta$ is its characteristic function then $\gamma_{j}(\Theta(0))=\gamma_{j}(T)$. Thus, let $T$ be a $C_{0}$ contraction with Jordan model $S\left(m_{1}\right) \oplus$ $\oplus S\left(m_{2}\right) \oplus \ldots$ such that $m_{1}(0) \neq 0$ and $\lim _{j \rightarrow \infty}\left|m_{j}(0)\right|=1$. Proposition 4.5 shows that the Jordan model of $T$ can be characterized within the class $\mathscr{T}$ of contractions,
which are quasisimilar with $T$ by its extremal properties. Indeed, define $\mathscr{T}_{n}$ recurrently, by $\mathscr{T}_{0}=\mathscr{T}$ and $\mathscr{T}_{n+1}=\left\{T^{\prime} \in \mathscr{T}_{n} \mid \gamma_{n+1}\left(T^{\prime}\right)=\inf _{s \in \mathscr{F}_{n}} \gamma_{n+1}(S)\right\}$. Then the only member up to unitary equivalence of $\bigcap_{n=0}^{\infty} \mathscr{T}_{n}$ is the Jordan model of $T$.

## Part II

## § 5. Preliminaries

1. We begin with a short review of the properties of infinite determinants (see [4], ch. IV, § 1), in order to discuss (in the next section) minors of such determinants.

Let $\Omega$ be a complex separable Hilbert space and $\mathscr{C}_{1}(\Omega)$ the ideal of nuclear operators, endowed with the trace-norm

$$
\begin{equation*}
\|X\|_{1}=\operatorname{tr}|X|,|X|=\left(X^{*} X\right)^{1 / 2} \quad\left(X \in \mathscr{C}_{1}(\Omega)\right) \tag{5.1}
\end{equation*}
$$

Consider $X \in I+\mathscr{C}_{1}(\Omega)$ and let $\left\{\lambda_{j}(X)\right\}_{j=1}^{\infty}$ be the eigenvalues of $X$ (repeated according to their multiplicities). We have

$$
\sum_{j=1}^{\infty}\left|1-\lambda_{j}(X)\right| \leqq \operatorname{tr}|I-X|<\infty
$$

and it follows that the infinite product defining the determinant

$$
\begin{equation*}
\operatorname{det}(X)=\prod_{j=1}^{\infty} \lambda_{j}(X) \tag{5.2}
\end{equation*}
$$

converges absolutely. Moreover, $\operatorname{det}(I+Y)$ as a function of $Y \in \mathscr{C}_{1}(\Omega)$ is analytic (in particular continuous on the Banach space $\mathscr{C}_{1}(\Omega)$ ). This follows from [4], Ch. IV, Corollary 1.1 and property $8^{\circ}$ on p. 207, combined with Proposition 2 on p. 11 of [3].

Also for $\left\{e_{j}\right\}_{j=1}^{\infty}$ an orthonormal basis of $\Omega$ and $X \in I+\mathscr{C}_{1}(\Omega)$, we have

$$
\begin{equation*}
\operatorname{det}(X)=\lim _{N \rightarrow \infty} \operatorname{det}\left[\left\langle X e_{i}, e_{j}\right\rangle\right]_{1 \leqq i, j \leqq N} \tag{5.3}
\end{equation*}
$$

(cf. [4], property $2^{\circ}$ on p. 203).
Furthermore, for $X, X^{\prime} \in I+\mathscr{C}_{1}(\Omega)$ we have (cf. the proof of property $7^{\circ}$ on p. 206 of [4]):

$$
\begin{equation*}
\operatorname{det}\left(X X^{\prime}\right)=\operatorname{det}(X) \operatorname{det}\left(X^{\prime}\right) \tag{5.4}
\end{equation*}
$$

In view of (5.2) the following assertions are easily seen to be true: a) if $X \in I+$ $+\mathscr{C}_{1}(\Omega)$ is unitary then $|\operatorname{det}(X)|=1$; b) if $X \in I+\mathscr{C}_{1}(\Omega)$ is a contraction then
$|\operatorname{det}(X)| \leqq 1 ;$ c) $X \in I+\mathscr{C}_{1}(\Omega)$ is invertible if and only if $\operatorname{det}(X) \neq 0 ;$ d) the determinant is invariant under similarities.
2. For any Hilbert space $\Omega$ we shall indicate by " $\rightarrow$ " the weak convergence in $\Omega$ and in $\mathscr{L}(\Omega)$. In order to avoid antilinear mappings we shall consider the dual space $\Omega^{d}$. If $T \in \mathscr{L}(\Omega)$, the dual operator is denoted by $T^{d}\left(T^{d} \in \mathscr{L}\left(\Omega^{d}\right)\right)$. $\left(\Omega^{d}\right)^{d}$ can be identified in the usual way with $\mathcal{A}$.
3. For any Hilbert space $\Omega$ and $n \geqq 0$ we shall denote by $\mathfrak{\Omega}^{\wedge n}$ the $n$-th exterior power of $\Omega$. For $n=0$ this is just the complex field $\mathbf{C}$ and in general $\Omega^{\wedge n}$ coincides with $\Omega^{\tau_{n}}$ for $\tau_{n}=(1,1,1, \ldots, 1,0, \ldots), \iota\left(\tau_{n}\right)=\operatorname{dim} \Omega, \&\left(\tau_{n}\right)=n$ (cf. §1.4). $\Omega^{\wedge n}$ is generated by vectors of the form

$$
\begin{equation*}
k_{1} \wedge k_{2} \wedge \ldots \wedge k_{n}=(n!)^{-1 / 2} \sum_{\sigma \in \mathbb{E}_{n}} \varepsilon(\sigma) k_{\sigma(1)} \otimes \ldots \otimes k_{\sigma(n)}, \quad k_{j} \in \Omega \quad(1 \leqq j \leqq n) \tag{5.5}
\end{equation*}
$$

where $\varepsilon(\sigma)$ is the sign of the permutation $\sigma$.
The factor ( $n!)^{-1 / 2}$ has been chosen so that $\left\|e_{1} \wedge \ldots \wedge e_{n}\right\|=1$ for any orthonormal system $\left\{e_{1}, e_{2}, \ldots, e_{n}\right\}$.

For $n, m$ two positive integers there is a bilinear map

$$
\wedge: \mathfrak{\Omega}^{\wedge n} \times \mathfrak{\Omega}^{\wedge m} \rightarrow \mathfrak{\Omega}^{\wedge(m+n)}
$$

such that $\left(k_{1} \wedge k_{2} \wedge \ldots \wedge k_{n}\right) \wedge\left(k_{n+1} \wedge \ldots \wedge k_{n+m}\right)=k_{1} \wedge \ldots \wedge k_{n+m}$. For each $A \in \mathscr{L}(\Omega)$ we shall denote $\varrho_{\tau_{n}}(A)$ as $A^{\wedge n}$, so that

$$
\begin{equation*}
A^{\wedge n}\left(k_{1} \wedge \ldots \wedge k_{n}\right)=A k_{1} \wedge \ldots \wedge A k_{n} \tag{5.6}
\end{equation*}
$$

Let $\Omega$ now be a Hilbert space of finite dimension $n$. If $\left\{e_{1}, \ldots, e_{n}\right\}$ is an orthonormal basis of $\Omega$, we can define a bilinear form

$$
B: \boldsymbol{\Omega}^{\wedge k} \times \mathfrak{\Omega}^{\wedge(n-k)} \rightarrow \mathbf{C}
$$

by the formula

$$
\begin{equation*}
B(h, g)=\left\langle h \wedge g, e_{1} \wedge \ldots \wedge e_{n}\right\rangle \tag{5.7}
\end{equation*}
$$

Choosing in $\boldsymbol{\Omega}^{\wedge \boldsymbol{j}}$ the usual orthonormal basis

$$
\left\{e_{i_{1}} \wedge \ldots \wedge e_{i_{j}} \mid 1 \leqq i_{1}<i_{2}<\ldots<i_{j} \leqq n\right\}
$$

it is easy to see that the mapping

$$
\begin{equation*}
C: \Omega^{\wedge(n-k)} \rightarrow\left(\Omega^{\wedge k}\right)^{d} \tag{5.8}
\end{equation*}
$$

given by $C(g)(h)=B(h, g)$ for $g \in \Omega^{\wedge(n-k)}, h \in \Omega^{\wedge k}$ is a linear isometry. If $A \in \mathscr{L}(\Omega)$ we have

$$
\begin{equation*}
B\left(A^{\wedge k} h, A^{\wedge(n-k)} g\right)=\operatorname{det}(A) B(h, g) \tag{5.9}
\end{equation*}
$$

because $A^{\wedge n}=\operatorname{det}(A) I_{g} \wedge n$. Let us define

$$
\begin{equation*}
F=C A^{\wedge(n-k)} C^{-1} \in \mathscr{L}\left(\left(\Im^{\wedge k}\right)^{\delta}\right) \tag{5.10}
\end{equation*}
$$

and

$$
\begin{equation*}
A^{A d k}=F^{d} \in \mathscr{L}\left(\Omega^{\wedge k}\right) \tag{5.11}
\end{equation*}
$$

We have $B\left(A^{A d k} h, g\right)=C(g)\left(A^{A d k} h\right)=(F(C(g)))(h)=\left(C\left(A^{\wedge(n-k)} g\right)\right)(h)=$ $=B\left(h, A^{\wedge(n-k)} g\right)$ and since $C$ is isometric,

$$
\begin{equation*}
\left\|A^{\wedge(n-k)}\right\|=\|F\|=\left\|A^{A d k}\right\| \tag{5.12}
\end{equation*}
$$

Also, as $B$ is nondegenerate we have

$$
\begin{equation*}
A^{A d k} A^{\wedge k}=\operatorname{det}(A) I_{\Omega \wedge k} \tag{5.13}
\end{equation*}
$$

It is obvious by the definition of $A^{A d k}$ that

$$
\begin{equation*}
\left(A_{1} A_{2}\right)^{A d k}=A_{2}^{A d k} A_{1}^{A d k}, A_{1}, A_{2} \in \mathscr{L}(\Omega) \tag{5.14}
\end{equation*}
$$

and it can be shown that

$$
\begin{equation*}
\left(A^{*}\right)^{A d k}=\left(A^{A d k}\right)^{*} \tag{5.15}
\end{equation*}
$$

Moreover, for invertible $A$ we infer from (5.13) that

$$
\begin{equation*}
A^{\wedge k} A^{A d k}=\operatorname{det}(A) I_{\Im \wedge k} \tag{5.16}
\end{equation*}
$$

and by continuity it follows that (5.16) always holds.
For $\left\{f_{1}, f_{2}, \ldots, f_{k}\right\}$ an orthonormal system in $\Omega$ we shall show that

$$
\begin{equation*}
\left\langle A^{A d k}\left(f_{1} \wedge \ldots \wedge f_{k}\right), f_{1} \wedge \ldots \wedge f_{k}\right\rangle=\operatorname{det}(P+(I-P) A(I-P)) \tag{5.17}
\end{equation*}
$$

where $P$ denotes the orthogonal projection onto the linear span of $\left\{f_{1}, f_{2}, \ldots, f_{k}\right\}$. Completing the system $\left\{f_{1}, \ldots, f_{k}\right\}$ to an orthonormal basis $\left\{f_{1}, \ldots, f_{n}\right\}$, we have

$$
\begin{gathered}
\left\langle A^{A d k}\left(f_{1} \wedge \ldots \wedge f_{k}\right), f_{1} \wedge \ldots \wedge f_{k}\right\rangle= \\
=\left\langle\left(A^{A d k}\left(f_{1} \wedge \ldots \wedge f_{k}\right)\right) \wedge f_{k+1} \wedge \ldots \wedge f_{n}, f_{1} \wedge \ldots \wedge f_{n}\right\rangle= \\
=B\left(A^{A d k}\left(f_{1} \wedge \ldots \wedge f_{k}\right), f_{k+1} \wedge \ldots \wedge f_{n}\right) \cdot\left\langle f_{1} \wedge \ldots \wedge f_{n}, e_{1} \wedge \ldots \wedge e_{n}\right\rangle^{-1}= \\
=B\left(f_{1} \wedge \ldots \wedge f_{k}, A^{\wedge(n-k)}\left(f_{k+1} \wedge \ldots \wedge f_{n}\right)\right) \cdot\left\langle f_{1} \wedge \ldots \wedge f_{n}, e_{1} \wedge \ldots \wedge e_{n}\right\rangle^{-1}= \\
=\left\langle f_{1} \wedge \ldots \wedge f_{k} \wedge A^{\wedge(n-k)}\left(f_{k+1} \wedge \ldots \wedge f_{n}\right), f_{1} \wedge \ldots \wedge f_{n}\right\rangle= \\
=\left\langle(P+A(I-P))^{\wedge n}\left(f_{1} \wedge \ldots \wedge f_{n}\right), f_{1} \wedge \ldots \wedge f_{n}\right\rangle= \\
=\operatorname{det}(P+A(I-P))=\operatorname{det}(P+(I-P) A(I-P)) .
\end{gathered}
$$

Formulas (5.14), (5.16), (5.17) show that $A^{A d k}$ does not depend on the particular choice of the orthonormal basis $\left\{e_{1}, \ldots, e_{n}\right\}$.

Let us now suppose that $A$ is a positive operator with eigenvalues $\lambda_{1} \geqq \lambda_{2} \geqq \ldots \geqq \lambda_{n}$ and the corresponding eigenvectors $f_{1}, f_{2}, \ldots, f_{n}$. Then $A^{\wedge(n-k)}$ is positive with
eigenvalues

$$
\lambda_{i_{1}} \lambda_{i_{2}} \ldots \lambda_{i_{n-k}}\left(1 \leqq i_{1}<i_{2}<\ldots<i_{n-k} \leqq n\right)
$$

It follows that

$$
\begin{aligned}
\left\|A^{A d k}\right\| & =\left\|A^{\wedge(n-k)}\right\|=\lambda_{1} \ldots \lambda_{n-k} \leqq\left(1+\left|\lambda_{1}-1\right|\right)\left(1+\left|\lambda_{2}-1\right|\right) \ldots\left(1+\left|\lambda_{n-k}-1\right|\right) \leqq \\
& \leqq \exp \left(\left|\lambda_{1}-1\right|\right) \exp \left(\left|\lambda_{2}-1\right|\right) \ldots \exp \left(\left|\lambda_{n-k}-1\right|\right) \leqq \exp (\operatorname{Tr}|A-I|) .
\end{aligned}
$$

Now for any $T \in \mathscr{L}(\boldsymbol{\Re})$, we have

$$
\left\|T^{*} T-I\right\|_{1} \leqq\left(1+\|T-I\|_{1}\right)^{2}-1 \quad \text { and } \quad\left\|\left(T^{*} T\right)^{1 / 2}-I\right\|_{1} \leqq\left\|T^{*} T-I\right\|_{1}
$$

as can be seen by comparing the eigenvalues of these operators. Therefore,

$$
\left\|\left(T^{*} T\right)^{1 / 2}-I\right\|_{1} \leqq\left(1+\|T-I\|_{1}\right)^{2}-1
$$

In particular, for the polar decomposition $T=U A$ of $T\left(A=|T|=\left(T^{*} T\right)^{1 / 2}\right)$ it follows that:

$$
\begin{gather*}
\left\|T^{A d k}\right\|=\left\|A^{A d k} U^{A d k}\right\| \leqq\left\|A^{A d k}\right\| \leqq \exp (\operatorname{Tr}|A-I|) \leqq  \tag{5.18}\\
\leqq \exp \left(\left(1+\|T-I\|_{1}\right)^{2}-1\right) .
\end{gather*}
$$

## § 6. Infinite dimensional adjoints and minors

Let us now consider $\Omega$ an infinite dimensional Hilbert space and $A \in \mathscr{L}(\Omega)$ so that $\operatorname{rank}(I-A)<\infty$. From the preceding considerations we easily infer the existence of an operator $A^{A d k} \in \mathscr{L}\left(\Omega^{\wedge k}\right)$ satisfying

$$
\begin{gather*}
A^{A d k} A^{\wedge k}=A^{\wedge k} A^{A d k}=\operatorname{det}(A) I_{\Re \wedge k}  \tag{6.1}\\
\left\langle A^{A d k}\left(f_{1} \wedge \ldots \wedge f_{k}\right), f_{1} \wedge \ldots \wedge f_{k}\right\rangle=\operatorname{det}(P+(I-P) A(I-P)), \tag{6.2}
\end{gather*}
$$

for $P$ the orthogonal projection onto the linear span of the orthonormal system $\left\{f_{1}, \ldots, f_{k}\right\}$;

$$
\begin{equation*}
\left\|A^{A d k}\right\| \leqq \exp \left(\left(1+\|A-I\|_{1}\right)^{2}-1\right) . \tag{6.3}
\end{equation*}
$$

Also for $A_{1}, A_{2} \in \mathscr{L}(\Omega)$ with rank $\left(I-A_{j}\right)<\infty, j=1,2$, we have

$$
\begin{equation*}
\left(A_{1} A_{2}\right)^{A d k}=A_{2}^{A d k} A_{1}^{A d k} . \tag{6.4}
\end{equation*}
$$

Let $A \in \mathscr{L}(\Omega)$ now be such that $I-A \in \mathscr{C}_{1}(\Omega)$ and let $A_{n}$ be such that $\operatorname{rank}\left(I-A_{n}\right)<\infty$ and $\lim _{n \rightarrow \infty}\left\|A-A_{n}\right\|_{1}=0$.

Using the fact that the function

$$
\mathscr{C}_{1}(\Omega) \ni X \rightarrow \operatorname{det}(I+X)
$$

is continuous, it follows from (6.2-3) that the sequence $A_{n}^{\text {Adk }}$ converges weakly. The limit, which will be denoted by $A^{\text {Adk }}$, satisfies (6.2-3). Because $A_{n}^{\wedge k}$ converges to $A^{\wedge k}$ in norm and $\operatorname{det}\left(A_{n}\right) \rightarrow \operatorname{det}(A)$ we also obtain property (6.1) for $A^{A d k}$. Using now (6.2-3) it follows that:

$$
\begin{equation*}
A, A_{n} \in I+\mathscr{C}_{1}(\Omega) \text { and }\left\|A_{n}-A\right\|_{1} \rightarrow 0 \quad \text { imply } A_{n}^{A d k} \rightarrow A^{A d k} \tag{6.5}
\end{equation*}
$$

Property (6.4) for $A_{1}, A_{2} \in I+\mathscr{C}_{1}(\Omega)$ follows from (6.1), provided $A_{1}, A_{2}$ are invertible, and can be extended using (6.5) to the case when only $A_{1}$ is invertible. Using (6.5) once again it follows that (6.4) holds in the general case.

We have shown in §5.1 that the function $Y \rightarrow \operatorname{det}(I+Y)$ is analytic on the Banach space $\mathscr{C}_{1}(\Omega)$. Using (6.2-3) we infer that for $\xi, \eta \in \mathfrak{\Omega}^{\wedge k}$ the mapping

$$
\mathbf{C} \ni \lambda \rightarrow\left\langle(I+X+\lambda Y)^{A d k} \xi, \eta\right\rangle
$$

is analytic when $X, Y \in \mathscr{C}_{1}(\Omega)$.
From this fact and from (6.3), using [3], Proposition 2 it follows that

$$
\mathscr{C}_{1}(\Omega) \ni X \rightarrow\left\langle(I+X)^{A d k} \xi, \eta\right\rangle
$$

for $\xi, \eta \in \boldsymbol{\Omega}^{\wedge k}$ is analytic.
This again implies the following stronger fact: the mapping

$$
\begin{equation*}
\mathscr{C}_{1}(\Omega) \ni X \rightarrow(I+X)^{A d k} \in \mathscr{L}\left(\Omega^{\wedge k}\right) \tag{6.6}
\end{equation*}
$$

is analytic (in particular continuous with respect to the norm topologies).
Let us also remark that for any contraction $A \in I+\mathscr{C}_{1}(\Omega)$ the adjoints $A^{A d k}$ are contractions. This is obvious if $\operatorname{dim} \Re=n<\infty$ (since in this case $\left\|A^{A d k}\right\|=$ $\left.=\left\|A^{\wedge(n-k)}\right\|\right)$ and follows in the general case by a simple limit argument.

We are now going to define the minors of an infinite determinant. Let $\mathfrak{M}$ and $\mathfrak{R}$ be two closed subspaces of $\Omega, P_{\mathfrak{M}}$ and $P_{\mathfrak{\Re}}$ the corresponding projections, and suppose there is a unitary operator $U \in I+\mathscr{C}_{1}(\Omega)$ such that $U \mathfrak{M}=\mathfrak{N}$. Then for $A \in I+\mathscr{C}_{1}(\Omega)$ the minor of $\operatorname{det}(A)$ corresponding to the triple $(\mathfrak{M}, \mathfrak{N}, U)$ is.

$$
\begin{equation*}
\operatorname{det}\left(U P_{\mathfrak{N}} A \mid \mathfrak{N}\right) \tag{6.7}
\end{equation*}
$$

The definition makes sense because it is easily seen that $U P_{\mathfrak{M}} A \mid \mathfrak{N} \in I_{\mathfrak{R}}+\mathscr{C}_{1}(\mathfrak{N})$. In case $\mathfrak{N}$ (and hence $\mathfrak{M}$ also) is of finite codimension in $\mathfrak{N}$, we shall say that $\operatorname{det}\left(U P_{\mathfrak{M}} A \mid \mathfrak{N}\right)$ is a minor of corank $\operatorname{dim} \mathfrak{M}^{\perp}$.

Let $\operatorname{det}\left(U P_{\mathfrak{M}} A \mid \mathfrak{N}\right)$ be a minor of corank $k$ of $A$. Then, by (6.2)

$$
\begin{gather*}
\operatorname{det}\left(U P_{\mathfrak{M}} A \mid \mathfrak{M}\right)=\operatorname{det}\left(P_{\mathfrak{\Re}} U A P_{\mathfrak{N}}+\left(I-P_{\mathfrak{Y}}\right)\right)=  \tag{6.8}\\
=\left\langle(U A)^{A d k}\left(e_{1} \wedge \ldots \wedge e_{k}\right), e_{1} \wedge \ldots \wedge e_{k}\right\rangle
\end{gather*}
$$

for $\left\{e_{1}, \ldots, e_{k}\right\}$ an orthonormal basis of $\Omega \ominus \mathfrak{N}$. Thus the minors of corank $k$ of $A$ coincide with some matrix elements of $(U A)^{A d k}=A^{A d k} U^{A d k}$.

## § 7. Determinants of contractive analytic functions

Let $\Theta \in H^{\infty}(\mathscr{L}(\Omega))$ be a contractive function (here $\Omega$ denotes as usual a separable Hilbert space). Let us suppose that $I-\Theta(\lambda)$ is nuclear for $\lambda \in D$ and let $\left\{e_{n}\right\}_{n=1}^{\infty}$ be an orthonormal basis of $\Omega$. The functions

$$
d_{n}(\lambda)=\operatorname{det}\left[\left\langle\Theta(\lambda) e_{i}, e_{j}\right\rangle\right]_{1 \leqq i, j \leqq n}=\operatorname{det}\left(P_{n} \Theta(\lambda) P_{n}+\left(1-P_{n}\right)\right)
$$

(here $P_{n}$ denotes the orthogonal projection onto the linear span of $\left\{e_{1}, \ldots, e_{n}\right\}$ ) are analytic,

$$
\begin{equation*}
\left|d_{n}(\lambda)\right| \leqq 1 \tag{7.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{n \rightarrow \infty} d_{n}(\lambda)=\operatorname{det}(\Theta(\lambda)) \tag{7.2}
\end{equation*}
$$

From (7.1) and (7.2) we infer, by the Vitali-Montel theorem, that $\operatorname{det}(\Theta(\lambda))$ is an analytic function. A similar argument shows that the functions $\lambda \rightarrow(\Theta(\lambda))^{\text {Adk }}$ are analytic and contractive (cf. §6) and that

$$
\begin{equation*}
\Theta^{\wedge k} \Theta^{A d k}=\Theta^{A d k} \Theta^{\wedge k}=\operatorname{det}(\Theta) I_{\boldsymbol{\Omega} \wedge k} \tag{7.3}
\end{equation*}
$$

In particular, if $\Theta(\lambda)$ is invertible for some $\lambda \in D$, it follows that $\Theta$ has a scalar multiple (cf. [13], ch. V, §6).

In case $\mathfrak{M}, \mathfrak{N}$ are subspaces of $\mathfrak{\Omega}$ of finite codimension and $U \in I+\mathscr{C}_{1}(\Omega)$ is a unitary operator such that $U \mathfrak{M}=\mathfrak{M}$, the function $\lambda \rightarrow \operatorname{det}\left(U P_{\mathfrak{M}} \Theta(\lambda) \mid \mathfrak{R}\right)$ is analytic and of modulus $\leqq 1$. We call such a function a minor of $\Theta$ of corank $\operatorname{dim} \mathfrak{M}^{\perp}$.

Let us denote by $\delta_{r}(\Theta)$ the greatest common inner divisor of the minors of corank $r$ of $\Theta(r=0,1,2, \ldots)$. For $r=0, \delta_{0}(\Theta)$ coincides with the inner factor of $\operatorname{det}(\Theta(\lambda))$. From (6.8) it follows that $\delta_{r}(\Theta)$ coincides with the greatest common inner divisor of the matrix elements of $\Theta^{\text {Adr }}$.

Lemma 7.1. $\delta_{r+1}(\Theta)$ divides $\delta_{r}(\Theta)$ for each $r$.
Proof. We have to prove that $\delta_{r+1}(\Theta)$ divides each minor of corank $r$ of $\Theta$. Clearly it suffices to prove that $\delta_{1}(\Theta)$ divides $\operatorname{det}(\Theta)$ or, equivalently,

$$
\operatorname{det}(\Theta) H^{2}(\Omega) \subset \delta_{1}(\Theta) H^{2}(\Omega)
$$

But this easily follows from the relation $\Theta \Theta^{A d 1}=\operatorname{det}(\Theta) I_{\Omega}$. Indeed, $\Theta^{A d 1} H^{2}(\Omega) \subset$ $\subset \delta_{1}(\Theta) H^{2}(\Omega)$ and, since $\Theta$ is analytic,

$$
\operatorname{det}(\Theta) H^{2}(\Omega)=\Theta \Theta^{A d 1} H^{2}(\Omega) \subset \Theta \delta_{1}(\Theta) H^{2}(\Omega) \subset \delta_{1}(\Theta) H^{2}(\Omega)
$$

Lemma 7.2. The greatest common inner divisor of the functions $\delta_{j}(\Theta)(j=1,2, \ldots)$ is 1 .

Proof. Let us denote by $m$ the greatest common inner divisor of the family $\left\{\delta_{j}(\Theta)\right\}_{0}^{\infty}$ and let $\left\{e_{j}\right\}_{j=1}^{\infty}$, be an orthonormal basis of $\Omega$. Since $\Theta^{A d r} H^{2}\left(\Omega^{\wedge r}\right) \subset$ $\subset m H^{2}\left(\boldsymbol{\Omega}^{\wedge r}\right)$ for each $r$, we have

$$
\begin{gathered}
|m(0)| \geqq\left|\left\langle\Theta(0)^{A d r}\left(e_{1} \wedge e_{2} \wedge \ldots \wedge e_{r}\right), e_{1} \wedge e_{2} \wedge \ldots \wedge e_{r}\right\rangle\right|= \\
=\left|\operatorname{det}\left(\left(I-P_{r}\right) \Theta(0)\left(I-P_{r}\right)+P_{r}\right)\right|,
\end{gathered}
$$

where $P_{r}$ denotes as usual the orthogonal projection onto the linear span of $\left\{e_{1}, \ldots, e_{r}\right\}$. We infer

$$
|m(0)| \geqq \limsup _{r \rightarrow \infty}\left|\operatorname{det}\left(\left(I-P_{r}\right) \Theta(0)\left(I-P_{r}\right)+P_{r}\right)\right|=1
$$

and the lemma follows.
Let us also note the relations

$$
\begin{equation*}
\delta_{j}\left(\Theta^{\sim}\right)=\delta_{j}(\Theta)^{\sim} \quad(j=1,2, \ldots) \tag{7.4}
\end{equation*}
$$

which hold for each function $\Theta$ of the type considered in this section.

## § 8. Weak contractions

Let us recall that a contraction $T$ acting on a Hilbert space $\mathfrak{S}$ is a weak contraction if its spectrum does not cover the unit disk $D$ and $I-T^{*} T$ is a nuclear operator. $T$ is a weak contraction if and only if $T^{*}$ is a weak contraction.

If a weak contraction $T$ is of class $C_{00}$ (that is $T^{n} \rightarrow 0$ and $T^{*^{n}} \rightarrow 0$ strongly as $n \rightarrow \infty$ ), then $T$ is of class $C_{0}$ and acts on a necessarily separable space. The proof of this fact goes as follows (cf. [13], Ch. VIII, § 1).

If we put

$$
\begin{equation*}
T_{\lambda}=(T-\lambda I)(I-\bar{\lambda} T)^{-1}, \quad \lambda \in D \tag{8.1}
\end{equation*}
$$

we have

$$
\begin{equation*}
I-T_{\lambda}^{*} T_{\lambda}=X_{\lambda}^{*}\left(I-T^{*} T\right) X_{\lambda}, X_{\lambda}=\left(1-|\lambda|^{2}\right)^{1 / 2}(I-\bar{\lambda} T)^{-1} \tag{8.2}
\end{equation*}
$$

So $T$ is a weak contraction if and only if $T_{\lambda}$ is a weak contraction. Moreover, we have $\left(T_{\lambda}\right)_{-\lambda}=T$. Therefore we may suppose without loss of generality that $T$ is invertible. Let $\left\{\mu_{j}\right\}_{1}^{n}\left(n \leqq \aleph_{0}\right)$ be the eigenvalues of $\left(I-T^{*} T\right) \mid \mathfrak{D}_{T}, \mathfrak{D}_{T}=\left(\left(I-T^{*} T\right) H\right)^{-}$ (multiple eigenvalues repeated according to their multiplicities). We have $\mu_{j} \neq 1$ because ker $T=\{0\}$.

Let $\left\{\varphi_{j}\right\}_{1}^{n}$ be an orthonormal basis of $\mathfrak{D}_{T}$ such that $\left(I-T^{*} T\right) \varphi_{j}=\mu_{j} \varphi_{j}$. It is easy to verify that the system $\left\{\psi_{j}\right\}_{1}^{n}$, where $\psi_{j}=\left(1-\mu_{j}\right)^{-1 / 2} T \varphi_{j}$, is an orthonormal basis of $\mathfrak{D}_{T^{*}}$ and that we have also $T^{*} \psi_{j}=\left(1-\mu_{j}\right)^{1 / 2} \varphi_{j}$.

Let us denote by $U$ the unitary operator determined by

$$
\begin{equation*}
U: \mathfrak{D}_{T} \rightarrow \mathfrak{D}_{T^{*}}, U \dot{\varphi}_{j}=-\psi_{j} \tag{8.3}
\end{equation*}
$$

Then the operator $(U+T) \mathfrak{D}_{T}$ is nuclear. Indeed,

$$
(U+T) h=\sum_{j=1}^{n}\left(\left(1-\mu_{j}\right)^{1 / 2}-1\right)\left(h, \varphi_{j}\right) \psi_{j}, \quad h \in \mathfrak{D}_{T}
$$

and from the relations
we infer

$$
\lim _{\mu \rightarrow 0} \mu^{-1}\left(1-(1-\mu)^{1 / 2}\right)=1 / 2, \quad \sum_{j=1}^{n} \mu_{j}<\infty
$$

$$
\sum_{j=1}^{n}\left(1-\left(1-\mu_{j}\right)^{1 / 2}\right)<\infty
$$

Furthermore, if $\Theta_{T} \in H_{i}^{\infty}\left(\mathscr{L}\left(\mathfrak{D}_{T}, \mathfrak{D}_{T^{*}}\right)\right)$ is the characteristic function of $T, U-\Theta_{T}(\lambda)$ is nuclear for $\lambda \in D$. Indeed,

$$
U-\Theta_{T}(\lambda)=(U+T)\left|\mathfrak{D}_{T}-\lambda D_{T^{*}}\left(I-\lambda T^{*}\right)^{-1} D_{T}\right| \mathfrak{D}_{T}\left(D_{T}=\left(I-T^{*} T\right)^{1 / 2}\right)
$$

and since $D_{T}$ and $D_{T^{*}}$ are Hilbert-Schmidt operators because $T$ is a weak contraction, $\lambda D_{T^{*}}\left(I-\lambda T^{*}\right)^{-1} D_{T}$ is nuclear. Thus the function $\Theta \in H_{i}^{\infty}\left(\mathscr{L}\left(\mathcal{D}_{T}\right)\right)$ defined by $\Theta(\lambda)=U^{*} \Theta_{T}(\lambda)$ coincides with $\Theta_{T}$ and $I-\Theta(\lambda)$ is nuclear for $\lambda \in D$.

Let us put

$$
\begin{equation*}
d_{T}(\lambda)=\operatorname{det}(\Theta(\lambda)), \quad \delta_{j}(T)=\delta_{j}(\Theta), \quad(j=0,1,2, \ldots) \tag{8.4}
\end{equation*}
$$

We have $d_{\mathrm{T}}(0)=\prod_{j=1}^{n}\left(1-\mu_{j}\right)^{1 / 2} \neq 0$ and from (7.3) (with $k=1$ ) it follows that $d_{T}$ is a scalar multiple of $\Theta$. As in [13], Theorem VI. 5.2 we obtain

Lemma 8.1. Each weak contraction $T$ of class $C_{00}$ is a $C_{0}$ contraction and its minimal function coincides with $\delta_{0}(T) / \delta_{1}(T)$.

Let us remark that we have a converse: suppose $\Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$ is such that $\Theta(\lambda) \in I+\mathscr{C}_{1}(\Omega), \lambda \in D$, and $\operatorname{det}(\Theta) \not \equiv 0$. Since $\operatorname{det}(\Theta)$ is then a scalar multiple of $\Theta$ (by (7.3) with $k=1$ ), it follows that $\Theta$ coincides with the characteristic function of an operator $T$ of class $C_{0}$ and from [13], Ch. IV § 1 it follows that $\operatorname{tr}\left(I-T^{*} T\right)=$ $=\operatorname{tr}\left(I-\Theta(0)^{*} \Theta(0)\right)<\infty$ so that $T$ is a weak contraction. Let us also note that the relations

$$
\begin{equation*}
d_{T^{*}}=d_{T}^{\sim}, \quad \delta_{j}\left(T^{*}\right)=\delta_{j}(T)^{\sim} \quad(j=0,1, \ldots) \tag{8.5}
\end{equation*}
$$

hold for each weak contraction $T$.
Proposition 8.2. Let $T$ be a weak $C_{0}$ contraction acting on the Hilbert space 5 and let $T=\left[\begin{array}{cc}T_{1} & X \\ 0 & T_{2}\end{array}\right], \mathfrak{H}^{\prime}=\mathfrak{H}_{1} \oplus \mathfrak{S}_{2}$ be the triangularization associated with the $T$-invariant subspace $\mathfrak{S}_{1}$. Then $T_{1}$ and $T_{2}$ are weak $C_{0}$ contractions and we have

$$
d_{T}=d_{T_{1}} d_{T_{2}}, \delta_{0}(T)=\delta_{0}\left(T_{1}\right) \delta_{0}\left(T_{2}\right)
$$

Proof. We may suppose without loss of generality that $T$ is invertible, thus $m_{T}(0) \neq 0$. By [13], Proposition III. 6.1, $T_{1}$ and $T_{2}$ are $C_{0}$ operators and $m_{T_{1}}, m_{T_{2}}$ are divisors of $m$. It follows that $m_{T_{1}}(0) \neq 0, m_{T_{2}}(0) \neq 0$ so that $T_{1}$ and $T_{2}$ are invertible. Moreover, we have

$$
I_{\mathfrak{5}_{1}}-T_{1}^{*} T_{1}=P_{\mathfrak{5}_{1}}\left(I-T^{*} T\right)\left|\mathfrak{S}_{1}, \quad I_{\mathfrak{5}_{2}}-T_{2} T_{2}^{*}=P_{\mathfrak{5}_{2}}\left(I-T T^{*}\right)\right| \mathfrak{S}_{2}
$$

thus $T_{1}$ and $T_{2}$ are weak contractions.
By [13] Theorem VII.1.1 and Proposition VII.2.1, we can associate with the invariant subspace $\mathfrak{S}_{1}$ a regular factorization

$$
\begin{equation*}
\Theta_{T}(\dot{\lambda})=\Theta_{2}(\lambda) \Theta_{1}(\lambda) \tag{8.6}
\end{equation*}
$$

such that the characteristic functions $\Theta_{T_{1}}(\lambda), \Theta_{T_{2}}(\lambda)$ coincide with the pure parts of $\Theta_{1}(\lambda), \Theta_{2}(\lambda)$, respectively. Then we have

$$
\Theta_{j}(\lambda)=U_{j}^{\prime}\left[\begin{array}{ll}
\Theta_{T_{j}}(\lambda) & 0  \tag{8.7}\\
0 & I_{j}
\end{array}\right] U_{j}^{\prime \prime}
$$

where $U_{j}^{\prime}, U_{j}^{\prime \prime}$ are unitary operators and $I_{j}$ denotes the identity operator on some Hilbert space ( $j=1,2$ ). Now, from the consideration preceding Lemma 8.1, it follows that $I-U_{j}^{0 *} \Theta_{T_{j}}(\lambda)$ is nuclear and $d_{T_{j}}(\lambda)=\operatorname{det}\left(U_{j}^{0 *} \Theta_{T_{j}}(\lambda)\right)$ for some unitary operators $U_{j}^{0}(j=1,2)$. With the notation

$$
U_{j}=U_{j}^{\prime}\left[\begin{array}{cc}
U_{j}^{0} & 0 \\
0 & I_{j}
\end{array}\right] U_{j}^{\prime \prime}
$$

we see that $I-U_{j}^{*} \Theta_{j}(\lambda)$ is nuclear and

$$
\begin{equation*}
d_{T_{j}}(\lambda)=\operatorname{det}\left(U_{j}^{*} \Theta_{j}(\lambda)\right) . \tag{8.8}
\end{equation*}
$$

Using (8.6) and (8.7) we obtain

$$
\begin{equation*}
U^{*} \Theta_{T}(\lambda)=U^{*} U_{2} U_{1}\left[U_{1}^{*}\left(U_{2}^{*} \Theta(\lambda)\right) U_{1}\right]\left(U_{1}^{*} \Theta_{1}(\lambda)\right) \tag{8.9}
\end{equation*}
$$

From this relation it follows that $I_{\mathfrak{D}_{r}}-U^{*} U_{2} U_{1}$ is a nuclear operator such that $\operatorname{det}\left(U^{*} U_{1} U_{2}\right)$ exists. Using (8.8-9) and (5.4) we then obtain

$$
\begin{aligned}
d_{T}(\lambda) & =\operatorname{det}\left(U^{*} U_{2} U_{1}\right) \operatorname{det}\left(U_{1}^{*}\left(U_{2}^{*} \Theta_{2}(\lambda)\right) U_{1}\right) \operatorname{det}\left(U_{1}^{*} \Theta_{1}(\lambda)\right)= \\
& =\operatorname{det}\left(U^{*} U_{2} U_{1}\right) \operatorname{det}\left(U_{2}^{*} \Theta_{2}(\lambda)\right) \operatorname{det}\left(U_{1}^{*} \Theta_{1}(\lambda)\right)= \\
& =\operatorname{det}\left(U^{*} U_{2} U_{1}\right) d_{T_{2}}(\lambda) d_{T_{1}}(\lambda) .
\end{aligned}
$$

The relation $\delta_{0}(T)=\delta_{0}\left(T_{1}\right) \delta_{0}\left(T_{2}\right)$ follows by taking the inner factors in the last obtained relations. The proposition is proved.

Remark 8.3. This proposition is a generalization of [13], Lemma IX. 3.1.

Lemma 8.4. A Jordan operator $S(M), M=\left\{m_{j}\right\}_{1}^{\infty}$, is a weak contraction if and only if $\sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|\right)<\infty$. In this case we have $d_{S(M)}=\delta_{0}(S(M))=\prod_{j=1}^{\infty} m_{j}$, where $\prod_{j=1}^{\infty} m_{j}$ means the limit of some converging subsequence of $\left\{m_{1} m_{2} \ldots m_{n}\right\}_{n=1}^{\infty}$.

Proof. For any inner function $m \in H^{\infty}$ we have

$$
\begin{aligned}
\left(I_{5(m)}-S(m) S(m)^{*}\right) & h=P_{5(m)}\left(I-U U^{*}\right) h=\left(h, c_{0}\right) P_{5(m)} c_{0}= \\
= & \left(h, c_{0}\right)(1-\overline{m(0)} m)
\end{aligned}
$$

( $h \in \mathfrak{G}(m)$ ), where $U$ denotes the unilateral shift on $H^{2}$ and $c_{0}$ is the constant functions $c_{0} \equiv 1$. Thus $I-S(m) S(m)^{*}$ is of rank one and has norm $\left(1-\overline{m(0)} m, c_{0}\right)=1-|m(0)|^{2}$. It follows that the trace norm of $I-S(M) S(M)^{*}$ equals $\sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|^{2}\right)$. We have only to remark that

$$
1-\left|m_{j}(0)\right| \leqq 1-\left|m_{j}(0)\right|^{2} \leqq 2\left(1-\left|m_{j}(0)\right|\right) .
$$

The equality $d_{S(M)}=\prod_{j=1}^{\infty} m_{j}$ obviously follows from the special form of the characteristic function of $S(M)$. So it remains only to prove that $\prod_{j=1}^{\infty} m_{j}$ is an inner function. To see this, let us remark that $\prod_{j=1}^{\infty} m_{j}$ and $\prod_{j=n^{n}}^{\infty} m_{j}$ have the same outer factor, such that this outer factor must be 1 because $\left|\prod_{j=n_{.}}^{\infty} m_{j}(-)\right| \rightarrow 1$ for each $\lambda \in D$. The lemma is proved.

From now on $T$ will denote a weak $C_{0}$ contraction acting on $\mathfrak{S}, \Theta \in H_{i}^{\infty}(\mathscr{L}(\Omega))$ will denote a function coinciding with the characteristic function of $T$ and $\Theta(\lambda) \in I+$ $+\mathscr{C}_{1}(\Omega), \lambda \in D$. We shall also denote by $S(M), M=\left\{m_{j}\right\}_{1}^{\infty}$, the Jordan model of $T$. From the relation

$$
\Theta^{\wedge r} \Theta^{A d r}=\Theta^{A d r} \Theta^{\wedge r}=d_{T} \cdot I_{\Omega \wedge r}, \quad \text { see }(7.3)
$$

we infer, because $\Theta^{\wedge r}$ is two-sided inner, that $\delta_{0}(T) / \delta_{r}(T)$ is the least inner scalar multiple of $\Theta^{\wedge r}$. Thus we have

$$
\begin{equation*}
d_{r}(T)=\delta_{0}(T) / \delta_{r}(T) \tag{8.10}
\end{equation*}
$$

Theorem 8.5. $A C_{0}$ contraction $T$ is a weak contraction if and only if its Jordan model $S(M), M=\left\{m_{j}\right\}_{1,}^{\infty}$, is a weak contraction.

Proof. That $T$ is a weak contraction if $S(M)$ is so follows from Proposition 4.3, via Lemma 8.4. So let us assume that $T$ is a weak contraction. Then, by Corollary 3.3.
and relation (8.10) it follows that $m_{1} m_{2} \ldots m_{r}$ divides $\delta_{0}(T)$ for each $r$. If we suppose $T$ is invertible, we have $\delta_{0}(T)(0) \neq 0$ and from the inequality

$$
\left|m_{1}(0) \ldots m_{r}(0)\right| \geqq\left|\delta_{0}(T)(0)\right|
$$

it follows that the infinite product $\prod_{j=1}^{\infty}\left|m_{j}(0)\right|$ converges. Therefore $\sum_{j=1}^{\infty}\left(1-\left|m_{j}(0)\right|\right)<\infty$ and our theorem follows by Lemma 8.4.

Proposition 8.6. For each weak $C_{0}$ contraction $T$, the determinant function $d_{T}$ is an inner function.

Proof. Let us write the inner-outer decomposition of $d_{T}$

$$
\begin{equation*}
d_{T}=d_{i} d_{0} \tag{8.11}
\end{equation*}
$$

Because $d_{i}$ is a scalar multiple of $\Theta^{\wedge k}$, there exists a contractive function $\Omega^{(k)} \in$ $\in H^{\infty}\left(\mathscr{L}\left(\boldsymbol{\Omega}^{\wedge k}\right)\right)$ such that

$$
\begin{equation*}
\Omega^{(k)} \Theta^{\wedge k}=\Theta^{\wedge k} \Omega^{(k)}=d_{i} I_{\mathrm{g} \wedge k} \tag{8.12}
\end{equation*}
$$

Then, by (7.3) and (8.12) we have

$$
\Theta^{k}\left(d_{0} \Omega^{(k)}-\Theta^{A d k}\right)=0
$$

so that ( $\Theta^{k}$ being inner)

$$
\begin{equation*}
\Theta^{A d k}=d_{o} \Omega^{(k)} \tag{8.13}
\end{equation*}
$$

Let $\left\{e_{i}\right\}_{i=1}^{\infty}$ be an orthonormal basis of $\Omega$ and denote by $P_{n}$ the orthogonal projection onto the linear span of $\left\{e_{1}, \ldots, e_{n}\right\}$. By (8.13) we have

$$
\left\langle\Theta^{A d k}\left(e_{1} \wedge \ldots \wedge e_{k}\right), e_{1} \wedge \ldots \wedge e_{k}\right\rangle=d_{o}\left\langle\Omega^{(k)}\left(e_{1} \wedge \ldots \wedge e_{k}\right), e_{1} \wedge \ldots \wedge e_{k}\right\rangle
$$

and therefore

$$
\begin{aligned}
& \left|d_{o}(0)\right| \geqq \limsup _{k \rightarrow \infty}\left|\left\langle\Theta(0)^{A d k}\left(e_{1} \wedge \ldots \wedge e_{k}\right), e_{1} \wedge \ldots \wedge e_{k}\right\rangle\right|= \\
& \quad=\limsup _{k \rightarrow \infty}\left|\operatorname{det}\left(\left(I-P_{k}\right) \Theta(0)\left(I-P_{k}\right)+P_{k}\right)\right|=1 .
\end{aligned}
$$

It follows that $\left|d_{o}(0)\right|=1$ so that $\left|d_{0}\right| \equiv 1$. The proposition follows.
We are now able to prove that the determinant function of a weak $C_{0}$ contraction is a quasi-similarity invariant.

Theorem 8.7. For each $C_{0}$ contraction $T$ with Jordan model $S(M), M=\left\{m_{j}\right\}_{1}^{\infty}$, we have

$$
\begin{align*}
& m_{j}=\delta_{j-1}(T) / \delta_{j}(T)  \tag{8.14}\\
& d_{T}=d_{S(M)}=\prod_{j=1}^{\infty} m_{j} \tag{8.15}
\end{align*}
$$

Proof. From (8.10) it follows that $\delta_{j-1}(T) / \delta_{j}(T)=d_{j}(T) / d_{j-1}(T)$ so the relation (8.14) obviously follows from Corollary 3.4.

For the second relation let us write (8.10) under the form

$$
\begin{equation*}
d_{\tau} \doteq \delta_{0}(T)=m_{1} m_{2} \ldots m_{n} \cdot \delta_{n}(T) \tag{8.16}
\end{equation*}
$$

(cf. Corollary 3.4). From Lemma 7.2 and Lemma 1 of [12] it follows that $d_{T}$ coincides with the least common inner multiple of the family $\left\{m_{1} m_{2} \ldots m_{n}\right\}_{n=1}^{\infty}$, which coincides with $d_{S(M)}$.

The theorem follows.

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| H. BERCOVICI | D. VOICULESCU |
| :--- | :--- |
| BD. BĂLCESCU 32 | INCREST |
| BUCURESTI, ROMANIA | CALEA VICTORIEI 114 |
|  | SECTOR 1 |
|  | BUCURESTI 22, ROMANIA |


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