## Techniques of Finsler geometry in the theory of vector bundles

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A topical problem in geometry is the study of the differential geometry of vector bundles. For this study the classical methods are not convenient enough, because of the very complicated analytical expressions of some important geometrical objects defined on the total space of a vector bundle.

Using the ideas from Finsler geometry [10, 11, 12, 13] we can considerably simplify the theory. For this, the notion of Finsler connection on the total space $E$ of a vector bundle $\xi=(E, \pi, M)$ is fundamental. To define it firstly we define the notion of the nonlinear connection $N$ on $E$, then we use this concept to obtain the algebra of Finsler tensor fields on $E$. A Finsler connection on the space $E$ is a linear connection $\nabla$ on $E$ which preserves by parallelism the horizontal distribution $N$ and the vertical distribution $E^{V}$ of $\xi$. $\nabla$ has four local coefficients which have very simple transformation laws to a change of canonical coordinates on $E$. Also, its torsion and its curvature have a small number of nontrivial components which are the Finsler tensor fields.

By applying this theory to the Riemannian structure $G$ on $E$ we get a Finsler canonical connection which has a simple form.

In this way the geometrical theory of $\xi$ can be constructed without difficulty. This method was used in our talk "Vector bundles Finsler geometry" presented at the second National Seminar on Finsler spaces [13], at the University of Braşov, Romania, February 15, 1982.

## 1. Vector bundles

Let $\xi=(E, \pi, M)$ be a vector bundle of the class $C^{\infty}$. We suppose that the total space $E$ has $n+m$ dimensions, the base $M$ has $n$ dimensions and the local fibre $E_{p}=\pi^{-1}(p), p \in M$, is a real vector space of dimension $n$.

If ( $U_{a}, \Phi_{a}$ ) is a vectorial chart of $\xi$, determined by the chart $\left(U_{\alpha}, \varphi_{a}\right)$ of the base manifold $M$, then $\Phi_{\alpha}: \pi^{-1}\left(U_{a}\right) \rightarrow U_{\alpha} \times \mathbf{R}^{m}$ is a diffeomorphism with the property $\mathrm{pr}_{1} \circ \Phi_{a}=\pi$ and $\Phi_{a, p}=\Phi_{a \mid E_{p}}: E_{p} \rightarrow \mathbf{R}^{m}$ is an isomorphism of vector spaces.

The mappings $g_{\beta \alpha}: U_{\beta} \cap U_{\alpha} \rightarrow G l(m, R)$, given by $g_{\beta \alpha}(p)=\Phi_{\beta, p} \circ \Phi_{\alpha, p}^{-1}, p \in U_{\beta} \cap$ $\cap U_{a}$, are the transition functions of $\xi$. To the vector chart ( $U_{a}, \Phi_{a}$ ) these corresponds a chart on the total space $E,\left(\pi^{-1}\left(U_{a}\right), \tilde{\Phi}_{a}\right)$, where $\tilde{\Phi}_{\alpha}$ is the diffeomorphism $\tilde{\Phi}_{a}=\left(\varphi_{a} \times 1_{\mathbf{R}^{m}}\right) \circ \Phi_{a}$. Therefore, the coordinates of the point $u \in \pi^{-1}\left(U_{a}\right) \subset E$ in this chart, for $u=\Phi_{a}^{-1}(p, y),(p, y) \in U_{a} \times \mathbf{R}^{m}$, are

$$
\widetilde{\Phi}_{a}(u)=\left(\varphi_{a} \times 1_{\mathbf{R}^{m}}\right) \circ \Phi_{\alpha} \circ \Phi_{a}^{-1}(p, y)=(x, y)=\left(x^{1}, \ldots, x^{n}, y^{1}, \ldots, y^{m}\right)
$$

and they are called the canonical coordinates of the point $u$ determinated by the coordinates $\left(x^{i}\right)$ of the point $p=\pi(u)$. Everywhere, the indices $i, j, k, l, \ldots$; $i^{\prime}, j^{\prime}, k^{\prime}, l^{\prime}, \ldots ; i^{\prime \prime}, j^{\prime \prime}, k^{\prime \prime}, l^{\prime \prime}, \ldots$ take the values $1,2, \ldots, n$ and $a, b, c, d, \ldots$; $a^{\prime}, b^{\prime}, c^{\prime}, d^{\prime}, \ldots ; a^{\prime \prime}, b^{\prime \prime}, c^{\prime \prime}, d^{\prime \prime}, \ldots$ take the values $1,2, \ldots, m$.

The transformations of the canonical coordinates $(x, y) \rightarrow\left(x^{\prime}, y^{\prime}\right)$ of the points of $E$, are given by:

$$
\left(x^{\prime}, y^{\prime}\right)=\tilde{\Phi}_{\beta} \circ \tilde{\Phi}_{\alpha}^{-1}(x, y)=\left(\left(\varphi_{\beta} \circ \varphi_{a}^{-1}\right)(x), g_{\beta x}(p) y\right)
$$

We write these transformations in the form:

$$
\begin{equation*}
x^{i^{\prime}}=x^{i^{\prime}}\left(x^{1}, \ldots, x^{n}\right), \quad y^{a^{\prime}}=M_{a}^{a^{\prime}}(x) y^{a}, \quad \operatorname{det}\left(M_{a}^{a^{\prime}}(x)\right) \neq 0 \tag{1.1}
\end{equation*}
$$

and the inverse transformations:

$$
x^{i}=x^{i}\left(x^{1^{\prime}}, \ldots, x^{n^{\prime}}\right), \quad y^{a}=\tilde{M}_{a^{\prime}}^{a}\left(x^{\prime}\right) y^{a^{\prime}}, \quad \operatorname{det}\left(\tilde{M}_{a^{\prime}}^{a}\left(x^{\prime}\right)\right) \neq 0 .
$$

The map $\pi: E \rightarrow M$ induces the $\pi^{T}$-morphism of the corresponding tangent bundles $\pi^{T}: T(E) \rightarrow T(M)$. Then $V E=\operatorname{Ker} \pi^{T}$ is. a subbundle of $T(E)$ called the vertical bundle. VE defines a distribution

$$
E^{V}: u \in E \rightarrow E_{u}^{V},
$$

where $E_{u}^{V}$ is the fibre of $V E$ in the point $u \in E . E^{V}$ is called the vertical distribution of $\xi$. On the open set $\pi^{-1}\left(U_{\alpha}\right), \frac{\partial}{\partial y^{a}}, a=1, \ldots, m$, is a local basis of the vertical distribution $E^{V}$. Therefore $E^{V}$ is integrable.

Definition'1.1. A non-linear connection on the total space $E$ of $\xi$ is a differentiable distribution $N: u \in E \rightarrow N_{u} \subset E_{u}$, with the property

$$
\begin{equation*}
E_{u}=N_{u} \oplus E_{u}^{V} \tag{1.2}
\end{equation*}
$$

where $E_{u}$ is the tangent space in the point $u$ to the manifold $E$.
It follows:

Proposition 1.1. If the base $M$ of the vector bundle $\xi$ is paracompact, then, on $E$, there exist the non-linear connections.

On $\pi^{-1}\left(U_{a}\right)$ the distribution $N$ has a local basis of the form

$$
\begin{equation*}
\frac{\delta}{\delta x^{i}}=\frac{\partial}{\partial x^{i}}-N_{i}^{a}(x, y) \frac{\partial}{\partial y^{a}} ; \tag{1.3}
\end{equation*}
$$

$N_{i}^{a}(x, y)$ are called the coefficients of the non-linear connection $N$. Also, $N$ can be locally determined by the Pfaff system $\delta y^{a}=0$, where

$$
\begin{equation*}
\delta y^{a}=d y^{a}+N_{i}^{a}(x, y) d x^{i} \tag{1.4}
\end{equation*}
$$

For every vector field $X$ on $E$ there exists a unique decomposition

$$
\begin{equation*}
X=X^{H}+X^{V}, \quad X_{u}^{H} \in N_{u}, \quad X_{u}^{V} \in E_{u}^{V}, \quad \forall u \in E . \tag{1.5}
\end{equation*}
$$

$X^{H}$ is called the horizontal part and $X^{V}$ the vertical part of $X$.
If $\omega$ is a 1 -form field on $E$, then we can define the 1 -form field $\omega^{H}$ on $E$ by the condition:

$$
\omega^{H}(X)=\omega\left(X^{H}\right), \quad \forall X \in \mathscr{X}(E)
$$

$\omega^{\boldsymbol{H}}$ is called the horizontal component of $\omega$. Let $\omega^{V}=\omega-\omega^{H}$. Then $\omega$ has a unique decomposition

$$
\begin{equation*}
\omega=\omega^{H}+\omega^{V} \tag{1.6}
\end{equation*}
$$

$\omega^{V}$ is called the vertical component of $\omega$. We have $\omega^{H}\left(X^{V}\right)=\omega^{V}\left(X^{H}\right)=0, \forall X \in \mathscr{X}(E)$.
Let us observe that $\frac{\delta}{\delta x^{i}}, i=1, \ldots, n$, being $n$ horizontal independent fields and $\partial / \partial y^{a}, a=1, \ldots, m$, being $m$ vertical independent fields, $\left(\delta / \delta x^{i}, \partial / \partial y^{a}\right)$ provides a local basis of the module of the vector fields $\mathscr{X}(E)$ and $\left(d x^{i}, \delta y^{a}\right)$ is a local basis of the module of 1 -form fields on $E$. These bases are dual:

$$
\begin{equation*}
\left\langle\frac{\delta}{\delta x^{i}}, d x^{j}\right\rangle=\delta_{i}^{j},\left\langle\frac{\delta}{\delta x^{i}}, \delta y^{a}\right\rangle=0,\left\langle\frac{\partial}{\partial y^{a}}, d x^{j}\right\rangle=0,\left\langle\frac{\partial}{\partial y^{a}}, \delta y^{b}\right\rangle=\delta_{a}^{b} . \tag{1.7}
\end{equation*}
$$

## 2. Algebra of Finsler tensor fields, Finsler connections

Definition 2.1. A tensor field $t$ on the total space $E$ of the vector bundle $\xi$ is called a Finsler tensor field of the type $\left(\begin{array}{ll}p & r \\ q & s\end{array}\right)$ if it has the property

$$
\begin{align*}
& t(\underset{1}{\omega}, \ldots, \underset{p}{\omega}, \underset{1}{X}, \ldots, \underset{q}{X}, \underset{p+1}{\omega}, \ldots, \underset{p+r}{\omega}, \underset{q+1}{X}, \ldots, \underset{q+s}{X})= \\
& =t\left(\omega_{1}^{H}, \ldots, \underset{p}{\omega^{H}}, \underset{1}{\boldsymbol{X}^{H}}, \ldots, \underset{q}{X^{H}}, \underset{p+1}{\omega}, \ldots, \underset{p+r}{\omega_{q+1}^{V}}, X_{q}^{V}, \ldots, X_{q+s}^{V}\right),  \tag{2.1}\\
& \underset{\alpha}{\forall \omega \in \mathscr{X}^{*}}(E), \quad \underset{\beta}{\forall X \in \mathscr{X}}(E) .
\end{align*}
$$

Proposition 2.1. A Finsler tensor field of the type $\left(\begin{array}{ll}p & r \\ q & s\end{array}\right)$ on $E$ has the following local form:

$$
\begin{equation*}
t=t_{j_{1} \ldots j_{q}, \ldots, i_{q}, b_{5}}^{\left.i_{1}, \ldots, y\right)}(x, y) \frac{\delta}{\delta x_{1}} \otimes \ldots \otimes \frac{\delta}{\delta x^{i_{p}}} \otimes d x^{j_{1}} \otimes \ldots \otimes d x^{j_{q}} \otimes \tag{2.2}
\end{equation*}
$$

$$
\otimes \frac{\partial}{\partial y^{a_{1}}} \otimes \ldots \otimes \frac{\partial}{\partial y^{a_{r}}} \otimes \delta y^{b_{1}} \otimes \ldots \otimes \delta y^{b_{e}} .
$$

The coordinate transformation (1.1) changes the coefficients $t_{j_{1} \ldots f_{b_{1}}{ }^{i} \ldots \ldots b_{s}}^{a_{1} \ldots a_{1}}(x, y)$ according to the law:

Theorem 2.1. If $w$ is a tensor field on $E$ of the type $(p, q)$ then it determines $2^{p+q}$ Finsler tensor fields on $E$ of the type $\left(\begin{array}{cc}p-r & r \\ q-s & s\end{array}\right)(r=0,1, \ldots, p ; s=0,1, \ldots, q)$.

Proof. The sum

$$
w\left(\underset { 1 } { \omega } \left(\underset{1}{\omega^{H}}+\underset{p}{\omega^{V}}, \ldots, \underset{p}{\omega^{H}}+\underset{1}{\omega^{V}}, \underset{1}{X^{H}}+\underset{1}{X^{V}}, \ldots, \underset{q}{X^{H}}+\underset{X^{V}}{X^{V}}\right.\right.
$$

has $2^{p+q}$ terms, each of them being a Finsler tensor field of the type mentioned.
The vector field $X^{H}$ is a Finsler tensor field of the type $\left(\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right)$, the vector field $X^{V}$ is a Finsler tensor field of the type $\left(\begin{array}{ll}0 & 1 \\ 0 & 0\end{array}\right)$, and $\omega^{H}, \omega^{V}$ are Finsler tensor fields of the type $\left(\begin{array}{ll}0 & 0 \\ 1 & 0\end{array}\right),\left(\begin{array}{ll}0 & 0 \\ 0 & 1\end{array}\right)$, respectively.

A remarkable local vector field is $y=y^{a} \frac{\partial}{\partial y^{a}}$. It is called the intrinsic vector field of $\xi$.
 then the $\mathscr{F}(E)$-module

$$
\underset{F}{\mathscr{T}}(E)=\underset{p, q, q, s=0,1, \ldots .}{ } \oplus_{\boldsymbol{F}}^{\mathscr{T}_{q}^{p r}}(E)
$$

and the product tensor is a graded algebra, called the algebra of Finsler tensor fields on $E$.

Definition 2.2. A Finsler connection on $E$ is a linear connection $\nabla$ on $E$ with the property that the horizontal linear spaces $N_{u}, u \in E$, of the distribution $N$ are parallel with respect to $\nabla$ and similarly, the vertical linear spaces $E_{u}^{V}, u \in E$, are parallel with respect to $\nabla$.

In what follows we shall prove the existence of a Finsler connection on $E$. We observe that a linear connection $\nabla$ on $E$ is a Finsler connection on $E$ if and only if

$$
\begin{equation*}
\left(\nabla_{X} Y^{H}\right)^{V}=0, \quad\left(\nabla_{X} Y^{V}\right)^{\boldsymbol{H}}=0, \quad \forall X, Y \in \mathscr{X}(E) \tag{2.3}
\end{equation*}
$$

Then, we have:
Theorem 2.2. The following statements are equivalent:
(a) $\nabla$ is a Finsler connection on $E$,
(b) $\nabla_{X} Y=\left(\nabla_{X} Y^{H}\right)^{H}+\left(\nabla_{X} Y^{V}\right)^{V}, \quad \forall X, Y \in \mathscr{X}(E)$,
(c) $\nabla_{X} \omega=\left(\nabla_{X} \omega^{H}\right)^{H}+\left(\nabla_{X} \omega^{V}\right)^{V}, \quad \forall \omega \in \mathscr{X}{ }^{*}(E), \quad \forall X \in \mathscr{X}(E)$.

For a Finsler connection $\nabla$ we put:

$$
\begin{equation*}
\nabla_{X}^{H} Y=\nabla_{X^{H}} Y, \quad \nabla_{X}^{V} Y=\nabla_{X^{V}} Y, \quad \forall X, Y \in \mathscr{X}(E) . \tag{2.4}
\end{equation*}
$$

The following theorem is easy to prove:
Theorem 2.3. For any Finsler connection $\nabla$ on $E, \nabla^{H}$ and $\nabla^{V}$ given by (2.4) are the covariant derivatives in the algebra $\underset{F}{\mathscr{T}}(E)$.
$\nabla^{H}$ is called the $h$-covariant derivative and $\nabla^{V}$ is called the $v$-covariant derivative of the Finsler connection $\nabla$.

Proposition 2.2. We have:
(1) $\nabla_{X}^{H} f=X^{H} f, \quad\left(\nabla_{X}^{H} Y^{H}\right)^{V}=0, \quad\left(\nabla_{X}^{H} Y^{V}\right)^{H}=0$,
(2) $\nabla_{X}^{H} Y=\left(\nabla_{X}^{H} Y^{H}\right)^{H}+\left(\nabla_{X}^{H} Y^{V}\right)^{V}$,
(3) $\nabla_{X}^{V} f=X^{V} f, \quad\left(\nabla_{X}^{V} Y^{H}\right)^{V}=0, \quad\left(\nabla_{X}^{V} Y^{V}\right)^{H}=0$,
(4) $\nabla_{X}^{V} Y=\left(\nabla_{X}^{V} Y^{H}\right)^{H}+\left(\nabla_{X}^{V} Y^{V}\right)^{V}$.

We have analogous formulas for $\nabla_{x} \omega$, too.
Theorem 2.4. If $\nabla^{H}$ and $\nabla^{V}$ are two covariant derivatives in the algebra of Finsler tensor fields $\underset{F}{\mathscr{F}}(E)$, having the properties (1) and (3) from Proposition 2.2, then there exists an unique Finsler connection $\nabla$ on $E$ for which $\nabla^{H}$ is the h-covariant derivative and $\nabla^{V}$ is the $v$-covariant derivative of $\nabla$.

If $t$ is a Finsler tensor field of the type $\left(\begin{array}{ll}p & r \\ q & s\end{array}\right)$, then for any Finsler connection $\nabla$ the following formulas hold:

$$
\begin{aligned}
& \nabla_{X} t=\nabla_{X}^{H} t+\nabla_{X}^{V} t, \\
& \left.\left(\nabla_{X}^{H} t\right) \underset{1}{(\omega, \ldots, \underset{q+s}{X})}=X^{H} t \underset{1}{\omega}, \ldots, \underset{q+s}{X}\right)-t\left(\nabla_{X}^{H} \underset{1}{\omega}, \ldots, \underset{q+s}{X}\right)-\ldots-t\left(\underset{1}{\omega}, \ldots, \nabla_{X}^{H} \underset{q+s}{X}\right), \\
& \left(\nabla_{X}^{V} t\right)(\underset{1}{\omega}, \ldots, \underset{q+s}{X})=X^{V} t(\underset{1}{\omega}, \ldots, \underset{q+s}{X})-t\left(\nabla_{X}^{V} \omega_{i}, \ldots, \underset{q+s}{X}\right)-\ldots-t\left(\underset{1}{\omega}, \ldots, \nabla_{X}^{V} \underset{q+a}{X}\right) .
\end{aligned}
$$

In. the canonical coordinates $\left(x^{i}, y^{a}\right)$ there exists a well-determined set of differentiable functions on $\pi^{-1}(U),\left(F_{j k}^{i}(x, y), F_{b k}^{a}(x, y), C_{j a}^{i}(x, y), C_{b c}^{a}(x, y)\right)$ such that

$$
\begin{cases}\nabla_{\partial \delta \delta x^{k}}^{H} \frac{\delta}{\delta x^{j}}=F_{j k}^{i}(x, y) \frac{\delta}{\delta x^{i}}, & \nabla_{\delta / \delta x^{k}}^{H} \frac{\partial}{\partial y^{b}}=F_{b k}^{a}(x, y) \frac{\partial}{\partial y^{a}}  \tag{2.5}\\ \nabla_{\partial / \partial y^{c}}^{V} \frac{\delta}{\delta x^{j}}=C_{j c}^{i}(x, y) \frac{\delta}{\delta x^{i}}, & \nabla_{\partial / \partial y^{c}}^{V} \frac{\partial}{\partial y^{b}}=C_{b c}^{a}(x, y) \frac{\partial}{\partial y^{a}}\end{cases}
$$

Let us consider, for simplicity, the Finsler tensor field

$$
K=K_{j b}^{i a}(x, y) \frac{\delta}{\delta x^{i}} \otimes \frac{\partial}{\partial y^{a}} \otimes d x^{j} \otimes \delta y^{b}
$$

Then, if we put for $X^{H}=X^{k} \frac{\delta}{\delta x^{k}}, X^{V}=X^{c} \frac{\partial}{\partial y^{c}}$,

$$
\nabla_{X}^{H} K=X^{k} K_{j b \mid k}^{i a} \frac{\dot{\delta}}{\delta x^{i}} \otimes \frac{\partial}{\partial y^{a}} \otimes d x^{j} \otimes \delta y^{b}, \quad \nabla_{X}^{V} K=\left.X^{c} K_{j b}^{i a}\right|_{c} \frac{\delta}{\delta x^{i}} \otimes \frac{\partial}{\partial y^{a}} \otimes d x^{j} \otimes \delta y^{b}
$$

we have:

$$
\begin{gather*}
K_{j b \mid k}^{i a}=\frac{\delta K_{j b}^{i a}}{\delta x^{k}}+F_{h k}^{i} K_{j b}^{h a}-F_{j k}^{h} K_{h b}^{i a}+F_{c k}^{a} K_{j b}^{i c}-F_{b k}^{c} K_{j c}^{i a}  \tag{2.6}\\
\left.K_{j b}^{i a}\right|_{c}=\frac{\partial K_{j b}^{i a}}{\partial y^{c}}+C_{h c}^{i} K_{j b}^{h a}-C_{j c}^{h} K_{h b}^{i a}+C_{d c}^{a} K_{j b}^{i d}-C_{b c}^{d} K_{j d}^{i a},
\end{gather*}
$$

which are the local expressions of the $h$ - and $v$-covariant derivatives.
Proposition 2.3. To a transformation of the canonical coordinates (1.1) the coefficients ( $F_{j k}^{i}, F_{b k}^{a}$ ) from (2.5) have the following transformation laws:

$$
\begin{gathered}
F_{j^{\prime} k^{\prime}}^{i^{\prime}}\left(x^{\prime} y^{\prime}\right)=\frac{\partial x^{i^{\prime}}}{\partial x^{i}} \frac{\partial x^{j}}{\partial x^{j^{\prime}}} \frac{\partial x^{k}}{\partial x^{k^{\prime}}} F_{j k}^{i}(x, y)+\frac{\partial x^{i^{\prime}}}{\partial x^{h}} \frac{\partial^{2} x^{h}}{\partial x^{j^{\prime}} \partial x^{k^{\prime}}}, \\
F_{b^{\prime} j^{\prime}}^{a^{\prime}}\left(x^{\prime}, y^{\prime}\right)=M_{a}^{a^{\prime}} \tilde{M}_{b^{\prime}}^{b} \frac{\partial x^{j}}{\partial x^{j^{\prime}}} F_{b j}^{a}(x, y)+M_{c}^{a^{\prime}} \frac{\partial \tilde{M}_{b^{\prime}}^{c}}{\partial x^{j^{\prime}}}
\end{gathered}
$$

and the coefficients $\left(C_{j c}^{i}, C_{b c}^{a}\right)$ from (2.5) are the Finsler tensor fields of the type $\left(\begin{array}{ll}1 & 0 \\ 1 & 1\end{array}\right),\left(\begin{array}{ll}0 & 1 \\ 0 & 2\end{array}\right)$, respectively.

Remark. If $N_{i}^{a}(x, y)$ are the coefficients of the non-linear connection $N$, then $\frac{\partial N_{k}^{a}}{\partial y^{b}}$ have the same transformation law as $F_{b k}^{a}(x, y)$.

These considerations allow to prove

Theorem 2.5. If $\xi=(E, \pi, M)$ is a vector bundle with the base $M$ paracompact, and $N$ is a fixed non-linear connection on the total space $E$, then there exist Finsler connections $\nabla$ on $E$, which preserve by parallelism the distributions $N$ and $E^{V}$.

The Finsler connection $\nabla$ with the coefficients $F_{b j}^{a}=\frac{\partial N_{j}^{a}}{\partial y^{b}}, C_{j b}^{i}=0, C_{b c}^{a}=0$ and $F_{j k}^{i}(x, y)$ arbitrary, is called a Berwald connection. Its simplicity made it very malleable in applications.

## 3. Torsion and curvature of Finsler connections

We consider again a non-linear connection $N$ on the total space $E$ of a vector bundle $\xi=(E, \pi, M)$ and let $\nabla$ be a Finsler connection on $E$, which preserves by parallelism the distributions $N$ and $E^{V}$.

Proposition 3.1. The torsion tensor field $T$ of the Finsler connection $\nabla$ is characterized by five Finsler tensor fields:
$\left[T\left(X^{H}, Y^{H}\right)\right]^{H}, \quad\left[T\left(X^{H}, Y^{H}\right)\right]^{V}, \quad\left[T\left(X^{H}, Y^{V}\right)\right]^{H}, \quad\left[T\left(X^{H}, Y^{V}\right)\right]^{V}, \quad\left[T\left(X^{V}, Y^{V}\right)\right]^{V}$.
If $T_{j k}^{i}, T_{j k}^{a}, P_{j b}^{i}, P_{j b}^{a}$ and $S_{b c}^{a}$ are their local components (where $T_{j k}^{i} \frac{\delta}{\delta x^{i}}=$ $=\left[T\left(\frac{\delta}{\delta x^{k}}, \frac{\delta}{\delta x^{j}}\right)\right]^{H}$, etc. $)$, then we have

$$
\begin{gather*}
T_{j k}^{i}=F_{j k}^{i}-F_{k j}^{i}, \quad T_{j k}^{a}=R_{j k}^{a}=\frac{\delta N_{j}^{a}}{\delta x^{k}}-\frac{\delta N_{k}^{a}}{\delta x^{j}}, \\
P_{j b}^{i}=C_{j b}^{i}, \quad P_{j b}^{a}=\frac{\partial N_{j}^{a}}{\partial y^{b}}-F_{b j}^{a}, \quad S_{b c}^{a}=C_{b c}^{a}-C_{c b}^{a} . \tag{3.2}
\end{gather*}
$$

It follows that the torsions $P_{j b}^{i}, P_{j b}^{a}, S_{b c}^{a}$ of a Berwald connection vanish.
Proposition 3.2. The Finsler tensor field $\left[T\left(X^{H}, Y^{H}\right)\right]^{V}$ vanishes if and only if the distribution $N$ is integrable.

The curvature tensor field $R$ of a Finsler connection $\nabla$ on $E$ has only six nontrivial Finsler components.

Proposition 3.3. The curvature tensor field $R$ of a Finsler connection $\nabla$ on the total space $E$ of $a$ vector bundle $\xi$ has the property

$$
\left[R(X, Y) Z^{H}\right]^{V}=\left[R(X, Y) Z^{V}\right]^{H}=0, \quad \forall X, Y, Z \in \mathscr{X}(E)
$$

Then, we have

Theorem 3.1. The curvature tensor field $R$ of a Finsler connection $\nabla$ on the total space $E$ of a vector bundle $\xi$ is characterized by the following six Finsler tensor fields:

$$
\begin{align*}
& \left\{\begin{array}{l}
R\left(X^{H}, Y^{H}\right) Z^{H}=\nabla_{X}^{H} \nabla_{Y}^{H} Z^{H}-\nabla_{Y}^{H} \nabla_{X}^{H} Z^{H}-\nabla_{\left[X^{H}, Y^{H}\right]}^{H} Z^{H}-\nabla_{\left[X^{H}, Y^{H}\right]}^{V} Z^{H}, \\
R\left(X^{H}, Y^{H}\right) Z^{V}=\nabla_{X}^{H} \nabla_{Y}^{H} Z^{V}-\nabla_{Y}^{H} \nabla_{X}^{H} Z^{V}-\nabla_{\left[X^{H}, Y^{H}\right]} Z^{V}-\nabla_{\left[X^{H}, Y^{H}\right]}^{V} Z^{V},
\end{array}\right.  \tag{3.3}\\
& \left\{\begin{array}{l}
R\left(X^{V}, Y^{H}\right) Z^{H}=\nabla_{X}^{V} \nabla_{Y}^{H} Z^{H}-\nabla_{Y}^{H} \nabla_{X}^{V} Z^{H}-\nabla_{\left[X^{V}, Y^{H}\right]}^{H} Z^{H}-\nabla_{\left[X^{V}, Y^{H}\right]}^{V} Z^{H}, \\
R\left(X^{V}, Y^{H}\right) Z^{V}=\nabla_{X}^{V} \nabla_{Y}^{H} Z^{V}-\nabla_{Y}^{H} \nabla_{X}^{V} Z^{V}-\nabla_{\left[X^{V}, Y^{H}\right]}^{H} Z^{V}-\nabla_{\left[X^{V}, Y^{H}\right]}^{V} Z^{V},
\end{array}\right.  \tag{3.3}\\
& \left\{\begin{array}{l}
R\left(X^{V}, Y^{V}\right) Z^{B}=\nabla_{X}^{V} \nabla_{Y}^{V} Z^{H}-\nabla_{Y}^{V} \nabla_{X}^{V} Z^{H}-\nabla_{\left[X^{V}, Y^{V}\right]}^{V} Z^{H}, \\
R\left(X^{V}, Y^{V}\right) Z^{V}=\nabla_{X}^{V} \nabla_{Y}^{V} Z^{V}-\nabla_{Y}^{V} \nabla_{X}^{V} Z^{V}-\nabla_{\left[X^{V}, Y^{V}\right]}^{V} Z^{V} .
\end{array}\right. \tag{3.3}
\end{align*}
$$

Let $R_{j k l}^{i}, R_{b k l}^{a}, P_{j k c}^{i}, P_{b k c}^{a}, S_{j b c}^{i}, S_{b c d}^{a}$ be the local components of the Finsler tensor fields (3.3), (3.3)', (3.3)", respectively $\left\{R\left(\frac{\delta}{\delta x^{l}}, \frac{\delta}{\delta x^{k}}\right) \frac{\delta}{\delta x^{j}}=R_{j k l}^{i} \frac{\delta}{\delta x^{i}}\right.$, etc. $\}$. Then, we have:

Theorem 3.2. The curvature tensor field $R$ of a Finsler connection $\nabla$ with the coefficients ( $F_{j k}^{i}, F_{b k}^{a}, C_{j c}^{i}, C_{b c}^{a}$ ) is characterized by the Finsler tensor fields (3.3), (3.3)', (3.3)", whose components are given by

$$
\begin{align*}
& \left\{\begin{array}{l}
R_{j k l}^{i}=\frac{\delta F_{j k}^{i}}{\delta x^{l}}-\frac{\delta F_{j l}^{i}}{\delta x^{k}}+F_{j k}^{h} F_{h l}^{i}-F_{j l}^{h} F_{h k}^{i}+C_{j a}^{i} R_{k l}^{a}, \\
R_{b k l}^{a}=\frac{\delta F_{b k}^{a}}{\delta x^{l}}-\frac{\delta F_{b l}^{a}}{\delta x^{k}}+F_{b k}^{c} F_{c l}^{a}-F_{b l}^{c} F_{c k}^{a}+C_{b c}^{a} R_{k l}^{c},
\end{array}\right.  \tag{3.4}\\
& \left\{\begin{array}{l}
P_{j k c}^{i}=\frac{\partial F_{j k}^{i}}{\partial y^{c}}-C_{j c \mid k}^{i}+C_{j b}^{i} P_{k c}^{b}, \\
P_{b k c}^{a}=\frac{\partial F_{b k}^{a}}{\partial y^{c}}-C_{b c \mid k}^{a}+C_{b d}^{a} P_{k c}^{d},
\end{array}\right.  \tag{3.4}\\
& \left\{\begin{array}{l}
S_{j b c}^{i}=\frac{\partial C_{j b}^{i}}{\partial y^{c}}-\frac{\partial C_{j c}^{i}}{\partial y^{b}}+C_{j b}^{h} C_{h c}^{i}-C_{j c}^{h} C_{h b}^{i}, \\
S_{b c d}^{a}=\frac{\partial C_{b c}^{a}}{\partial y^{d}}-\frac{\partial C_{b d}^{a}}{\partial y^{c}}+C_{b c}^{e} C_{e d}^{a}-C_{b d}^{e} C_{c c}^{a} .
\end{array}\right.
\end{align*}
$$

Observe the simplicity of these expressions, as compared to the components of the curvature tensor field $R$ written in the natural frame $\left(\frac{\partial}{\partial x^{i}}, \frac{\partial}{\partial y^{a}}\right)$, [4].

The Ricci and Bianchi identities can be written down without difficulty.

## 4. Riemannian structures; Finsler connections compatible with a metrical structure

The efficiency of the techniques of Finsler geometry in the study of vector bundles is visible. In particular, its applications to the theory of several geometric structures on the total space of a vector bundles are useful.

We study the Riemannian structure $G$ on the total space $E$ of a vector bundle $\xi=(E, \pi, M)$. If $E_{u}^{V}$ is the vertical space at the point $u$ of $E$, then the vectors orthogonal to $E_{u}^{V}$, with respect to $G$, uniquely determine the vector subspace $N_{u}$ complementary to $E_{u}^{V}$. That is, $E_{u}=N_{u} \oplus E_{u}^{V}, u \in E$, and the map $N: u \rightarrow N_{u}$ defines, in a geometrical way, a non-linear connection $N$ on the total space $E$.

Proposition 4.1. If $G$ is a Riemannian structure on $E$, then there exists an unique non-linear connection $N$ on $E$ with the property

$$
\begin{equation*}
G(X, Y)=0, \quad \forall X \in N, \quad \forall Y \in E^{v} \tag{4.1}
\end{equation*}
$$

Proposition 4.2. For a Riemannian structure $G$ on $E$, there exist an unique symmetric Finsler tensor field $G^{H}$ of the type $\left(\begin{array}{ll}0 & 0 \\ 2 & 0\end{array}\right)$ non-degenerate on the fibres of the horizontal bundle HE, and an unique symmetric Finsler tensor field $G^{V}$ of the type $\left(\begin{array}{ll}0 & 0 \\ 0 & 2\end{array}\right)$ non-degenerate on the fibres of the vertical bundle VE, such that

$$
\begin{equation*}
G=G^{H}+G^{V} \tag{4.2}
\end{equation*}
$$

Indeed, from (4.1) we get $G(X, Y)=G\left(X^{H}, Y^{H}\right)+G\left(X^{V}, Y^{V}\right)$. Putting $G^{H}(X, Y)=G\left(X^{H}, Y^{H}\right)$ and $G^{V}(X, Y)=G\left(X^{V}, Y^{V}\right)$ one obtains two Finsler tensor fields with the properties mentioned above.

A Finsler connection $\nabla$ on $E$, which preserves by parallelism the distributions $N$ and $E^{V}$, is called compatible with the Riemannian structure $G$ or is called a metrical Finsler connection if $\nabla_{X} G=0, \forall X \in X(E)$.

Proposition 4.3. A Finsler connection $\nabla$ on the total space $E$ is a metrical connection, with respect to the Riemannian structure $G$, if and only if

$$
\begin{equation*}
\nabla_{X}^{H} G^{H}=0, \quad \nabla_{X}^{H} G^{V}=0, \quad \nabla_{X}^{V} G^{H}=0, \quad \nabla_{X}^{V} G^{V}=0 . \tag{4.3}
\end{equation*}
$$

Theorem 4.1. If $G$ is a Riemannian structure on the total space $E$ of a vector bundle $\xi=(E, \pi, M)$, then the following Finsler connection is compatible with the structure $G$ :

$$
\begin{aligned}
2 G^{H}\left(\nabla_{X}^{H} Y^{H},\right. & \left.Z^{H}\right)=X^{H} G^{H}\left(Y^{H}, Z^{H}\right)+Y^{H} G^{H}\left(Z^{H}, X^{H}\right)-Z^{H} G^{H}\left(X^{H}, Y^{H}\right)- \\
& -G^{H}\left(X^{H} ;\left[Y^{H}, Z^{H}\right]^{H}\right)+G^{H}\left(Y^{H},\left[Z^{H}, X^{H}\right]^{H}\right)+G^{H}\left(Z^{H},\left[X^{H}, Y^{H}\right]^{H}\right),
\end{aligned}
$$

$$
\begin{align*}
& \nabla_{X}^{H} Y^{V}=\dot{\nabla}_{X}^{H} Y^{V}+B\left(Y^{V}, X^{H}\right), G^{V}\left(B\left(Y^{V}, X^{H}\right), Z^{V}\right)=(1 / 2)\left(\dot{\nabla}_{X}^{H} G^{V}\right)\left(Y^{V}, Z^{V}\right) \\
& \nabla_{X}^{V} Y^{H}=\dot{\nabla}_{X}^{V} Y^{H}+D\left(Y^{H}, X^{V}\right), G^{\dot{H}}\left(D\left(Y^{H}, X^{V}\right), Z^{H}\right)=(1 / 2)\left(\dot{\nabla}_{X}^{V} G^{H}\left(Y^{H}, Z^{H}\right)\right.  \tag{4.4}\\
& 2 G^{V}\left(\nabla_{X}^{V} Y^{V}, Z^{V}\right)=X^{V} G^{V}\left(Y^{V}, Z^{V}\right)+Y^{V} G^{V}\left(Z^{V}, X^{V}\right)-Z^{V} G^{V}\left(X^{V}, Y^{V}\right)- \\
&-G^{V}\left(X^{V},\left[Y^{V}, Z^{V}\right]^{V}\right)+G^{V}\left(Y^{V},\left[Z^{V}, X^{V}\right]^{V}\right)+G^{V}\left(Z^{V},\left[X^{V}, Y^{V}\right]^{V}\right),
\end{align*}
$$

where $\dot{\vee}$ is a fixed Finsler connection, which preserves by parallelism the distributions $N$ and $E^{V}$.

Proof. We know that there is a Finsler connection $\dot{\nabla}$ which preserves by parallelism the distributions $N$ and $E^{V}$. In the condition $\left[T\left(X^{H}, Y^{H}\right)\right]^{H}=0$, by the classical method, the first equality in (4.4) gives, uniquely, $\nabla_{X}^{H} Y^{H}$ and the second one $\nabla_{X}^{H} Y^{V}$. It is easy to see that $\nabla^{H}$, determined in this way, is a $h$-covariant derivative in the algebra $\mathscr{F}(E)$ and that we have $\nabla_{X}^{H} G^{H}=0, \nabla_{X}^{H} G^{V}=0$. Analogously, the third equation in (4.4) gives uniquely $\nabla_{X}^{V} Y^{H}$, and the last one, in condition $\left[T\left(X^{V}, Y^{V}\right)\right]^{V}=0$, allows to determine uniquely $\nabla_{X}^{V} Y^{V}$. This $\nabla^{V}$ is a $v$-covariant derivative in the algebra $\mathscr{T}(E)$ and it has the properties $\nabla_{X}^{V} G^{H}=0, \nabla_{X}^{V} G^{V}=0$. Therefore $\nabla=\nabla^{H}+\nabla^{V}$ is a Finsler connection on $E$ compatible with the Riemannian structure $G$.

In canonical coordinates, let $g_{i j}=G\left(\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta x^{j}}\right), g_{a b}=G\left(\frac{\partial}{\partial y^{a}}, \frac{\partial}{\partial y^{b}}\right)$. The equations (4.4) give the following coefficients of the metrical Finsler connection $\nabla$ :

$$
\begin{align*}
& F_{j k}^{i}(x, y)=\frac{1}{2} g^{i h}\left(\frac{\delta g_{h j}}{\delta x^{k}}+\frac{\delta g_{h k}}{\delta x^{j}}-\frac{\delta g_{j k}}{\delta x^{h}}\right), \quad F_{b k}^{a}(x, y)=\dot{F}_{b k}^{a}(x, y)+\frac{1}{2} g^{a c}, g_{b c i k} \\
& C_{j b}^{i}(x, y)=\dot{C}_{j b}^{i}(x, y)+\frac{1}{2} g^{i h} g_{j h} \dot{l}_{b}, \quad C_{b c}^{a}(x, y)=\frac{1}{2} g^{a d}\left(\frac{\partial g_{d b}}{\partial y^{c}}+\frac{\partial g_{d c}}{\partial y^{b}}-\frac{\partial g_{b c}}{\partial y^{d}}\right), \tag{4.5}
\end{align*}
$$

where ( $\dot{F}_{j k}^{i}, \dot{F}_{b k}^{a}, \dot{C}_{j b}^{i}, \dot{C}_{b c}^{a}$ ) are the coefficients of a fixed Finsler connection $\stackrel{\rightharpoonup}{V}$, and $i, i$ are the $h$ - and $v$-covariant derivatives with respect to $\dot{\nabla}$.

Observing that the distribution $N$ is geometrically determined by the structure $G$ we get that $F_{j k}^{i}(x, y)$ and $C_{b c}^{a}(x, y)$ are well-determined by $G$. Then, considering as a fixed Finsler connection $\stackrel{\vee}{ }$ the Finsler connection with coefficients

$$
\begin{equation*}
\dot{F}_{j k}^{\dot{j}}(x, y)=F_{j k}^{j}(x, y), \stackrel{\circ}{F}_{b k}^{a}(x, y)=\frac{\partial N_{k}^{a}}{\partial y^{b}}, \quad \dot{C}_{j b}^{i}(x, y)=0, \quad \dot{C}_{b c}^{a}=C_{b c}^{a}, \tag{4.6}
\end{equation*}
$$

we have
Theorem 4.2. The Finsler connection (4.5), (4.6) is metrical and depends only on the Riemannian structure $G$.

This connection can be called the canonical metrical Finsler connection. We get without difficulty:

Theorem 4.3. The Riemann-Christoffel connection of a Riemannian structure $G$ on the total space $E$ of a vector bundle $\xi$ coincides with the canonical metrical Finsler connection of $G$ if and only if
(1) the horizontal distribution $N$ is integrable,
(2) the metrical tensor field $G$ is constant on the fibres of vertical subbundle $V E$,
(3) $\left[T\left(X^{H}, Y^{V}\right)\right]^{V}=0$.

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