

Local analytic solutions to some nonhomogeneous problems with p -Laplacian*

Gabriella Bognár
University of Miskolc
Miskolc-Egyetemváros, 3515 Hungary

Abstract

Applying the Briot-Bouquet theorem we show that there exists an unique analytic solution to the equation $(t^{n-1}\Phi_p(y'))' + (-1)^i t^{n-1}\Phi_q(y) = 0$, on $(0, a)$, where $\Phi_r(y) := |y|^{r-1}y$, $0 < r, p, q \in \mathbf{R}^+$, $i = 0, 1$, $1 \leq n \in \mathbf{N}$, a is a small positive real number. The initial conditions to be added to the equation are $y(0) = A \neq 0$, $y'(0) = 0$, for any real number A . We present a method how the solution can be expanded in a power series for near zero.

1 Preliminaries

We consider the quasilinear differential equation

$$\Delta_p u + (-1)^i |u|^{q-1} u = 0, \quad u = u(x), \quad x \in \mathbf{R}^n,$$

where $n \geq 1$, p and q are positive real numbers, $i = 0, 1$ and Δ_p denotes the p -Laplacian $(\Delta_p u = \operatorname{div}(|\nabla u|^{p-1} \nabla u))$. If $n = 1$, then the equation is reduced to

$$(\Phi_p(y'))' + (-1)^i \Phi_q(y) = 0,$$

where for $r \in \{p, q\}$

$$\Phi_r(y) := \begin{cases} |y|^{r-1} y, & \text{for } y \in \mathbf{R} \setminus \{0\} \\ 0, & \text{for } y = 0. \end{cases}$$

We note that function Φ_r is an odd function. For $n > 1$ we restrict our attention to radially symmetric solutions. The problem under consideration is reduced to

$$(t^{n-1}\Phi_p(y'))' + (-1)^i t^{n-1}\Phi_q(y) = 0, \quad \text{on } (0, a) \quad (1)$$

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where $a > 0$. A solution of (1) means a function $y \in C^1(0, a)$ for which $t^{n-1}\Phi_p(y') \in C^1(0, a)$ and (1) is satisfied. We shall consider the initial values

$$\begin{aligned} y(0) &= A \neq 0, \\ y'(0) &= 0, \end{aligned} \tag{2}$$

for any $A \in \mathbf{R}$.

For the existence and uniqueness of radial solutions to (1) we refer to [9]. If $n = 1$ and $i = 0$, then it was showed that the initial value problem (1) – (2) has a unique solution defined on the whole \mathbf{R} (see [3], and [4]), moreover, its solution can be given in closed form in terms of incomplete gamma functions [4]. If $n = 1$, $i = 0$, Lindqvist gives some properties of the solutions [8]. If $n = 1$ and $p = q = 1$, then (1) is a linear differential equation, and its solutions are well-known:

if $i = 0$, the solution (1) – (2) with $A = 1$ is the cosine function,

if $i = 1$, the solution (1) – (2) with $A = 1$ is the hyperbolic cosine function,

and both the cosine and hyperbolic cosine functions can be expanded in power series.

In the linear case, when $n = 2$, $p = q = 1$, $i = 0$, the solution of (1) – (2) with $A = 1$ is $J_0(t)$, the Bessel function of first kind with zero order, and for $n = 3$, $p = q = 1$, $i = 0$ then the solution of (1) – (2) with $A = 1$ is $j_0(t) = \sin t/t$, called the spherical Bessel function of first kind with zero order.

In the cases above, for special values of parameteres n , p , q , i , we know the solution in the form of power series.

The type of singularities of (1) – (2) was classified in [1] in the case when $i = 0$, and $p = q$. If $n = 1$, then a solution of (1) is not singular.

Our purpose is to show the existence of the solution of problem (1) – (2) in power series form near the origin. We intend to examine the local existence of an analytic solution to problem (1) – (2) and we give a constructive procedure for calculating solution y in power series near zero. Moreover we present some numerical experiments.

2 Existence of an unique solution

We will consider a system of certain differential equations, namely, the special Briot-Bouquet differential equations. For this type of differential equations we refer to the book of E. Hille [6] and E. L. Ince [7].

Theorem 1 (*Briot-Bouquet Theorem*) *Let us assume that for the system of equations*

$$\left. \begin{aligned} \xi \frac{dz_1}{d\xi} &= u_1(\xi, z_1(\xi), z_2(\xi)), \\ \xi \frac{dz_2}{d\xi} &= u_2(\xi, z_1(\xi), z_2(\xi)), \end{aligned} \right\} \tag{3}$$

where functions u_1 and u_2 are holomorphic functions of ξ , $z_1(\xi)$, and $z_2(\xi)$ near the origin, moreover $u_1(0, 0, 0) = u_2(0, 0, 0) = 0$, then a holomorphic solution

of (3) satisfying the initial conditions $z_1(0) = 0, z_2(0) = 0$ exists if none of the eigenvalues of the matrix

$$\begin{bmatrix} \left. \frac{\partial u_1}{\partial z_1} \right|_{(0,0,0)} & \left. \frac{\partial u_1}{\partial z_2} \right|_{(0,0,0)} \\ \left. \frac{\partial u_2}{\partial z_1} \right|_{(0,0,0)} & \left. \frac{\partial u_2}{\partial z_2} \right|_{(0,0,0)} \end{bmatrix}$$

is a positive integer.

For a proof of Theorem 1 we refer to [2].

The differential equation (1) has singularity around $t = 0$ for the case $n > 1$.

Theorem 1 ensures the existence of formal solutions $z_1 = \sum_{k=0}^{\infty} a_k \xi^k$ and $z_2 = \sum_{k=0}^{\infty} b_k \xi^k$ for system (3), and also the convergence of formal solutions.

We apply the method Parades and Uchiyama [10].

Theorem 2 For any $p \in (0, +\infty), q \in (0, +\infty), i = 0, 1, n \in \mathbf{N}$ the initial value problem (1) $y(0) = A, y'(0) = 0$ has an unique analytic solution of the form $y(t) = Q(t^{1+1/p})$ in $(0, a)$ for small real value of a , where Q is a holomorphic solution to

$$Q'' = \frac{(-1)^{i+1}}{p(1+1/p)^{p+1}} t^{-\frac{p+1}{p}} \frac{\Phi_q(Q)}{|Q'|^{p-1}} - \frac{n}{p\alpha} t^{-(1+1/p)} Q'$$

near zero satisfying $Q(0) = A, Q'(0) = \frac{p}{p+1} \Phi_{1/p} [(-1)^{i+1} \Phi_q(A)/n]$.

Proof. We shall now present a formulation of (1) as a system of Briot-Bouquet type differential equations (3). Let us take solution of (1) in the form

$$y(t) = Q(t^\alpha), \quad t \in (0, a),$$

where function $Q \in C^2(0, a)$ and α is a positive constant. Substituting $y(t) = Q(t^\alpha)$ into (1) we get that Q satisfies

$$Q''(t^\alpha) = \frac{(-1)^{i+1}}{p\alpha^{p+1}} t^{-(\alpha-1)(p+1)} \frac{\Phi_q(Q)}{|Q'|^{p-1}} - \frac{n-1+p(\alpha-1)}{p\alpha} t^{-\alpha} Q'$$

and introducing variable ξ by $\xi = t^\alpha$ we have

$$Q''(\xi) = \frac{(-1)^{i+1}}{p\alpha^{p+1}} \xi^{-\frac{(\alpha-1)(p+1)}{\alpha}} \frac{\Phi_q(Q)}{|Q'|^{p-1}} - \frac{n-1+p(\alpha-1)}{p\alpha} \xi^{-1} Q'. \quad (4)$$

Here, we introduce function Q as follows

$$Q(\xi) = \gamma_0 + \gamma_1 \xi + z(\xi), \quad (5)$$

where $z \in C^2(0, a)$, $z(0) = 0$, $z'(0) = 0$. Therefore Q has to fulfill the properties $Q(0) = \gamma_0$, $Q'(0) = \gamma_1$, $Q'(\xi) = \gamma_1 + z'(\xi)$, $Q''(\xi) = z''(\xi)$. From initial condition $y(0) = A$ we have that

$$\gamma_0 = A.$$

We restate (4) as a system of equations:

$$\left. \begin{array}{l} z_1(\xi) = z(\xi) \\ z_2(\xi) = z'(\xi) \end{array} \right\} \text{ with } \left. \begin{array}{l} z_1(0) = 0 \\ z_2(0) = 0 \end{array} \right\},$$

according to (4) we get that

$$\begin{aligned} z''(\xi) &= \frac{(-1)^{i+1}}{p \alpha^{p+1}} \xi^{-\frac{(\alpha-1)}{\alpha}(p+1)} \frac{\Phi_q(\gamma_0 + \gamma_1 \xi + z(\xi))}{|\gamma_1 + z'(\xi)|^{p-1}} \\ &\quad - \frac{n-1+p(\alpha-1)}{p \alpha} \xi^{-1} (\gamma_1 + z'(\xi)). \end{aligned}$$

We generate the system of equations

$$\left. \begin{array}{l} u_1(\xi, z_1(\xi), z_2(\xi)) = \xi z_1'(\xi) \\ u_2(\xi, z_1(\xi), z_2(\xi)) = \xi z_2'(\xi) \end{array} \right\}$$

as follows

$$\left. \begin{array}{l} u_1(\xi, z_1(\xi), z_2(\xi)) = \xi z_2 \\ u_2(\xi, z_1(\xi), z_2(\xi)) = \frac{(-1)^{i+1}}{p \alpha^{p+1}} \xi^{1-p(\frac{\alpha-1}{\alpha})} \frac{\Phi_q(\gamma_0 + \gamma_1 \xi + z_1(\xi))}{|\gamma_1 + z_2(\xi)|^{p-1}} \\ \quad - \frac{n-1+p(\alpha-1)}{p \alpha} (\gamma_1 + z_2(\xi)) \end{array} \right\}.$$

In order to satisfy conditions $u_1(0, 0, 0) = 0$ and $u_2(0, 0, 0) = 0$ we must get zero for the power of ξ in the right-hand side of the second equation:

$$\frac{1-p(\alpha-1)}{\alpha} = 0,$$

i.e., $\alpha = \frac{1}{p} + 1$. To ensure $u_2(0, 0, 0) = 0$ we have the connection

$$n \Phi_p(\gamma_1) + \left(\frac{p}{p+1}\right)^p (-1)^i \Phi_q(\gamma_0) = 0,$$

i.e.,

$$\gamma_1 = (-1)^{i+1} \frac{p}{p+1} \Phi_{1/p} \left((-1)^{i+1} \frac{\Phi_q(\gamma_0)}{n} \right). \quad (6)$$

Therefore, taking into consideration that Φ_r is an even function for any $r \in \{p, q\}$, we obtain

$$\gamma_1 = \begin{cases} \frac{p}{p+1} A^{q/p} (-1)^{i+1} \frac{1}{n^{1/p}} & \text{if } A > 0, \\ \frac{p}{p+1} |A|^{q/p} (-1)^i \frac{1}{n^{1/p}} & \text{if } A < 0. \end{cases} \quad (7)$$

From initial conditions $y(0) = A \neq 0$, $y'(0) = 0$, and (5) it follows that $\gamma_0 = A$.

For u_1 and u_2 we find that

$$\begin{aligned} \frac{\partial u_1}{\partial z_1} \Big|_{(0,0,0)} &= 0, & \frac{\partial u_1}{\partial z_2} \Big|_{(0,0,0)} &= 0, \\ \frac{\partial u_2}{\partial z_1} \Big|_{(0,0,0)} &= -\frac{p^p q |\gamma_0|^{q-1}}{(p+1)^{p+1} |\gamma_1|^{p-1}}, & \frac{\partial u_2}{\partial z_2} \Big|_{(0,0,0)} &= -\frac{np}{p+1}. \end{aligned}$$

Therefore the eigenvalues of matrix

$$\begin{bmatrix} \partial u_1 / \partial z_1 & \partial u_1 / \partial z_2 \\ \partial u_2 / \partial z_1 & \partial u_2 / \partial z_2 \end{bmatrix}$$

at $(0, 0, 0)$ are 0 and $-np/(p+1)$. Since both eigenvalues are non-positive, applying Theorem 1 we get the existence of unique analytic solutions z_1 and z_2 at zero. Thus we get the analytic solution $Q(\xi) = \gamma_0 + \gamma_1 \xi + z(\xi)$ satisfying (4) with $Q(0) = \gamma_0$, $Q'(0) = \gamma_1$, where $\gamma_0 = A$ and γ_1 is determined in (7). ■

Corollary 3 *From Theorem 2 it follows that solution $y(t)$ for (1) has an expansion near zero of the form $y(t) = \sum_{k=0}^{\infty} a_k t^{k(\frac{1}{p}+1)}$ satisfying $y(0) = A$ and $y'(0) = 0$.*

3 Determination of local solution

We give a method for the determination of power series solution of (1) – (2). For simplicity, we take $A = 1$. Thus initial conditions

$$\begin{aligned} y(0) &= 1, \\ y'(0) &= 0 \end{aligned}$$

are considered. We seek a solution of the form

$$y(t) = a_0 + a_1 t^{\frac{1}{p}+1} + a_2 t^{2(\frac{1}{p}+1)} + \dots, \quad t > 0, \quad (8)$$

with coefficients $a_k \in \mathbf{R}$, $k = 0, 1, \dots$. From Section 2 we get that $a_0 = \gamma_0 = 1$ and $a_1 = \gamma_1 = \frac{p}{p+1}(-1)^{i+1} \frac{1}{n^{1/p}}$. Near zero $y(t) > 0$ and $y'(t) < 0$ for $i = 0$, $y'(t) > 0$ for $i = 1$. Therefore

$$\Phi_q(y(t)) = y^q(t) = \left(a_0 + a_1 t^{\frac{1}{p}+1} + a_2 t^{2(\frac{1}{p}+1)} + \dots \right)^q.$$

After differentiating (8), we get

$$y'(t) = t^{\frac{1}{p}} \left[a_1 \left(\frac{1}{p} + 1 \right) + 2a_2 \left(\frac{1}{p} + 1 \right) t^{\frac{1}{p}+1} + 3a_3 \left(\frac{1}{p} + 1 \right) t^{2(\frac{1}{p}+1)} + \dots \right],$$

and hence

$$\begin{aligned} & \Phi_p(y'(t)) = (-1)^{i+1} (y'(t))^p \\ & = (-1)^{i+1} t \left[a_1 \left(\frac{1}{p} + 1 \right) + 2a_2 \left(\frac{1}{p} + 1 \right) t^{\frac{1}{p}+1} + 3a_3 \left(\frac{1}{p} + 1 \right) t^{2(\frac{1}{p}+1)} + \dots \right]^p. \end{aligned}$$

For $y^q(t)$ and $(y'(t))^p$

$$y^q(t) = A_0 + A_1 t^{\frac{1}{p}+1} + A_2 t^{2(\frac{1}{p}+1)} + \dots \quad (9)$$

$$(y'(t))^p = t \left[B_0 + B_1 t^{\frac{1}{p}+1} + B_2 t^{2(\frac{1}{p}+1)} + \dots \right], \quad (10)$$

where coefficients A_k and B_k can be expressed in terms of a_k ($k = 0, 1, \dots$).

Using (10) we obtain

$$\begin{aligned} & (t^{n-1} \Phi_p(y'))' = \left((-1)^{i+1} t^n \left[B_0 + B_1 t^{\frac{1}{p}+1} + B_2 t^{2(\frac{1}{p}+1)} + \dots \right] \right)' \\ & = (-1)^{i+1} t^{n-1} \left[B_0 n + B_1 \left(n + \frac{1}{p} + 1 \right) t^{\frac{1}{p}+1} + B_2 \left(n + 2 \left(\frac{1}{p} + 1 \right) \right) t^{2(\frac{1}{p}+1)} + \dots \right], \end{aligned}$$

and substituting it to the equation (1) with (9) we get

$$\begin{aligned} & (-1)^{i+1} t^{n-1} \left[B_0 n + B_1 \left(n + \frac{1}{p} + 1 \right) t^{\frac{1}{p}+1} + B_2 \left(n + 2 \left(\frac{1}{p} + 1 \right) \right) t^{2(\frac{1}{p}+1)} + \dots \right] \\ & + (-1)^i t^{n-1} \left[A_0 + A_1 t^{\frac{1}{p}+1} + A_2 t^{2(\frac{1}{p}+1)} + \dots \right] = 0. \end{aligned}$$

Comparing the coefficients of the proper power of t we find

$$\begin{aligned} B_0 n - A_0 &= 0, \\ B_1 \left(n + \frac{1}{p} + 1 \right) - A_1 &= 0, \\ B_2 \left(n + 2 \left(\frac{1}{p} + 1 \right) \right) - A_2 &= 0, \\ &\vdots \\ B_k \left(n + k \left(\frac{1}{p} + 1 \right) \right) - A_k &= 0, \\ &\vdots \end{aligned} \quad (11)$$

Applying the J. C. P. Miller formula (see [5]) for the determination of A_k and B_k ($k = 0, 1, \dots$) we have.

$$A_k = \frac{1}{k} \sum_{j=0}^{k-1} [(k-j)q - j] A_j a_{k-j}, \quad (12)$$

$$B_k = \frac{p}{a_1 k(p+1)} \sum_{j=0}^{k-1} [(k-j)p - j] B_j a_{k-j+1} \left[(k-j+1) \left(\frac{1}{p} + 1 \right) \right] \quad (13)$$

for any $k > 0$.

From initial condition $y(0) = 1$ we get $a_0 = 1$, $A_0 = 1$, and therefore

$$B_0 = \frac{1}{n}.$$

From (11) for $i = 1$ we get $B_1(n + \frac{1}{p} + 1) - A_1 = 0$, and evaluating A_1 from (12) and B_1 from (13) we find

$$B_0 = \left[a_1 \left(\frac{1}{p} + 1 \right) \right]^p,$$

thus

$$a_1 = \frac{p}{p+1} (-1)^{i+1} \frac{1}{n^{1/p}}.$$

Similarly, we determine coefficients a_k for all $k = 0, 1, \dots$ from (11), (12) and (13).

Example 4 Solve (1) – (2) for $n=2$; $i=0$; $p=0.5$; $q=1$.

The solution of the differential equation $(t\Phi_{0.5}(y'))' + t\Phi_1(y) = 0$ with conditions $y(0) = 0$, $y'(0) = 1$ near zero we evaluate by MAPLE from (11), (12) and (13). We obtain

$$y(t) = 1 - 0.2222222222t^3 + 0.0370370370t^6 - 0.0047031158t^9 + 0.0005443421t^{12} + \dots$$

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